

Recent progress of understanding 3D magnetic topology in stellarators and tokamaks

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1. Introduction
2. HINT: a 3D MHD equilibrium calculation code
3. Impacts of beta-sequences to magnetic topology in stellarators and tokamaks
4. Impacts of toroidal rotation on magnetic islands
5. Summary and outlook

1.Introduction

2.HINT: a 3D MHD equilibrium calculation code

3.Impacts of beta-sequences to magnetic topology in stellarators and tokamaks

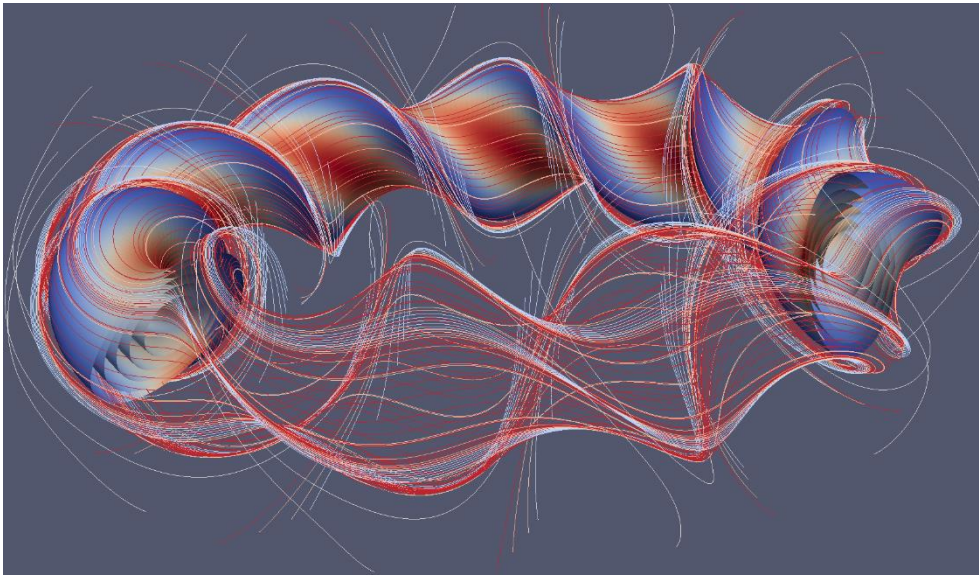
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5.Summary and outlook

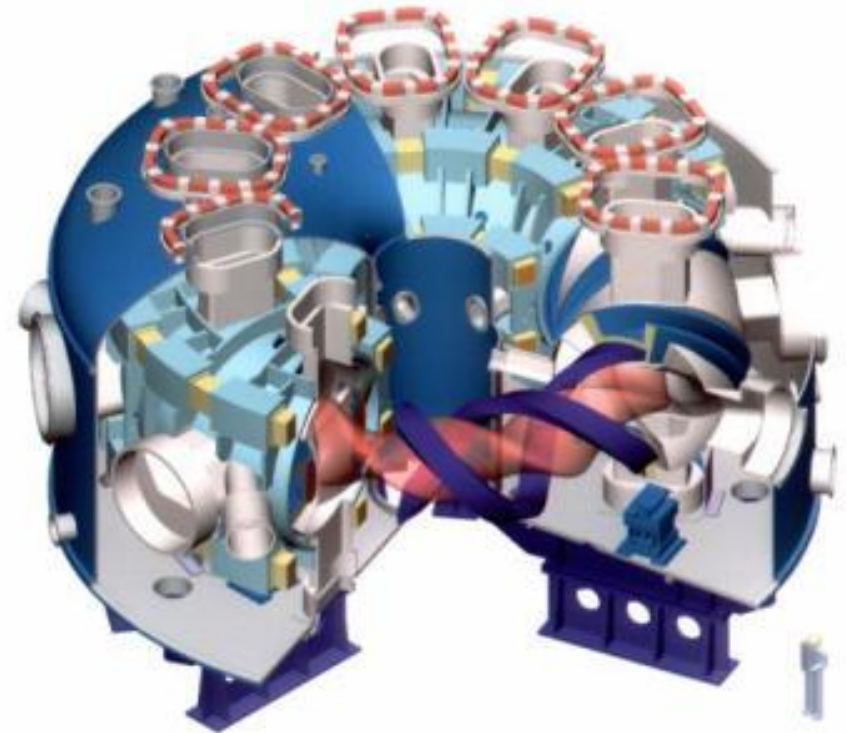
Background

- LHD is an $L/M=2/10$ Heliotron configuration. The stochastization of magnetic field lines naturally appears because of no symmetry.

Flux surfaces



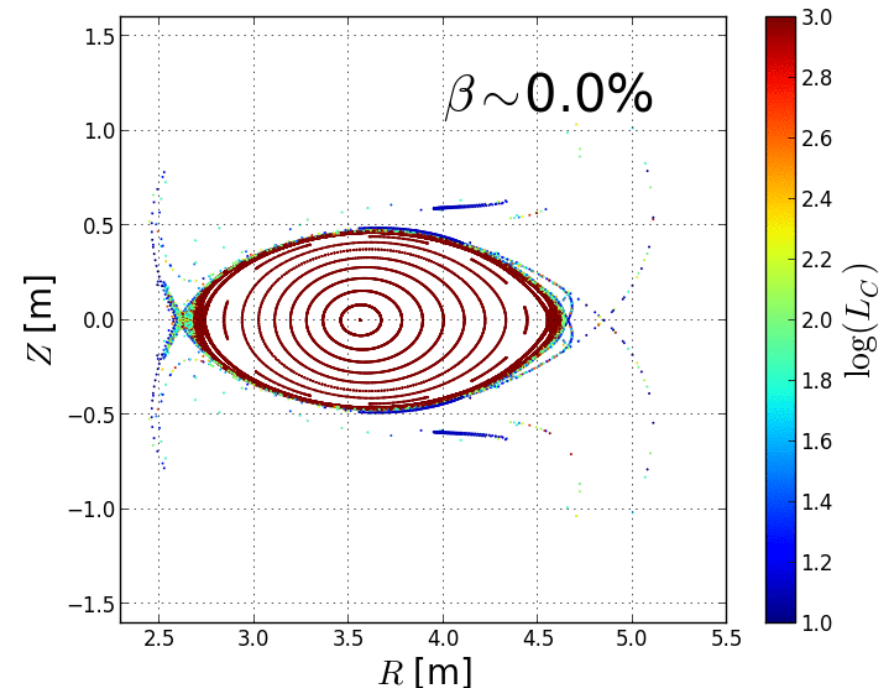
Large Helical Device



Background

- LHD is an L/M=2/10 Heliotron configuration. The stochastization of magnetic field lines naturally appears because of no symmetry.
- In addition, 3D MHD equilibrium analyses predict the stochastization by the “***nonlinear 3D equilibrium response***”. That is, magnetic field lines are further stochastized by pressure-induced perturbed field driven by currents along rippled field lines.

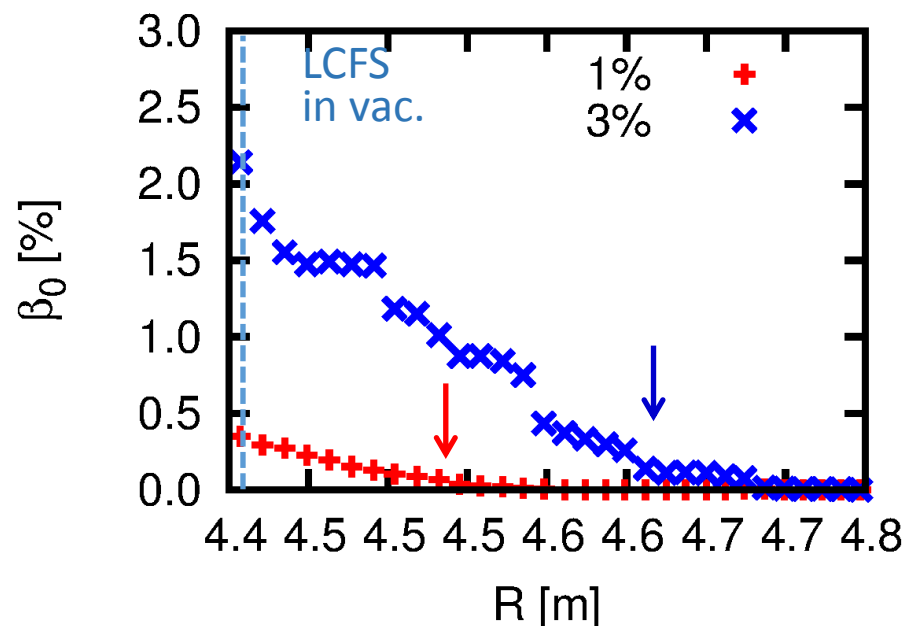
3D MHD modeling by HINT



Background

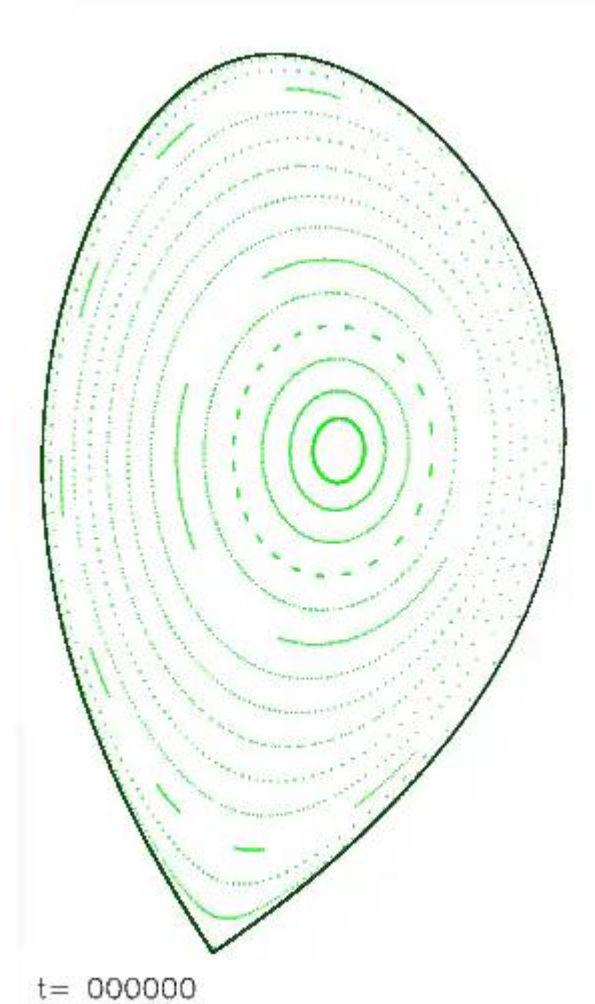
- LHD is an L/M=2/10 Heliotron configuration. The stochastization of magnetic field lines naturally appears because of no symmetry.
- In addition, 3D MHD equilibrium analyses predict the stochastization by the “***nonlinear 3D equilibrium response***”. That is, magnetic field lines are further stochastized by pressure-induced perturbed field driven by currents along rippled field lines.
- In experiments, changing the boundary of plasma pressure is observed. Increasing β , the boundary shifts to the outward of the torus.

Pressure profiles with different β



Magnetic topology changed!

2D boundary in axisymmetric torus

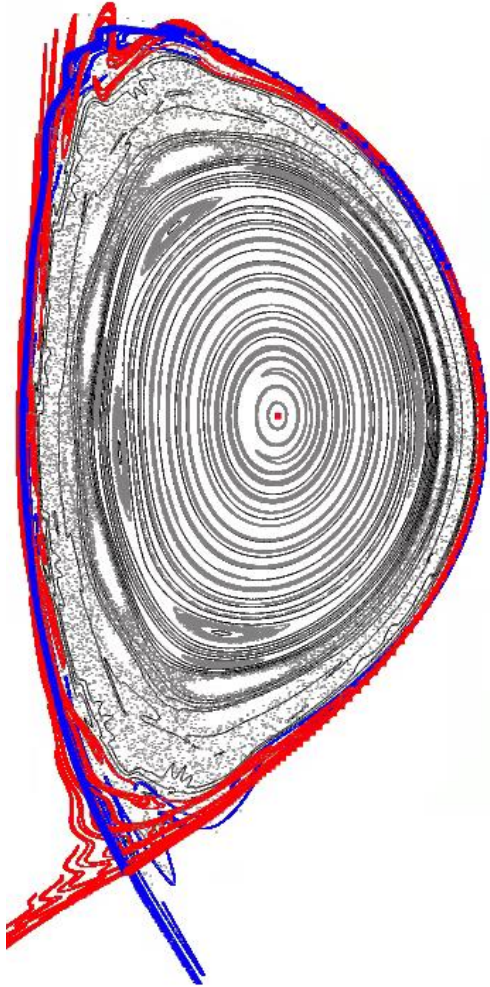


**In axisymmetric torus (2D system),
plasma boundary can be defined clearly
on separatrix or limiter.**



**2D plasma boundary separates nested flux surface
region and scrape-off layer (SOL), open field line
region.**

3D stochastic boundary in non-axisymmetric system



However,

In non-axisymmetric torus (general 3D system), perfectly nested flux surfaces cannot be guaranteed.



The magnetic field of 3D system consists of 3 regions.

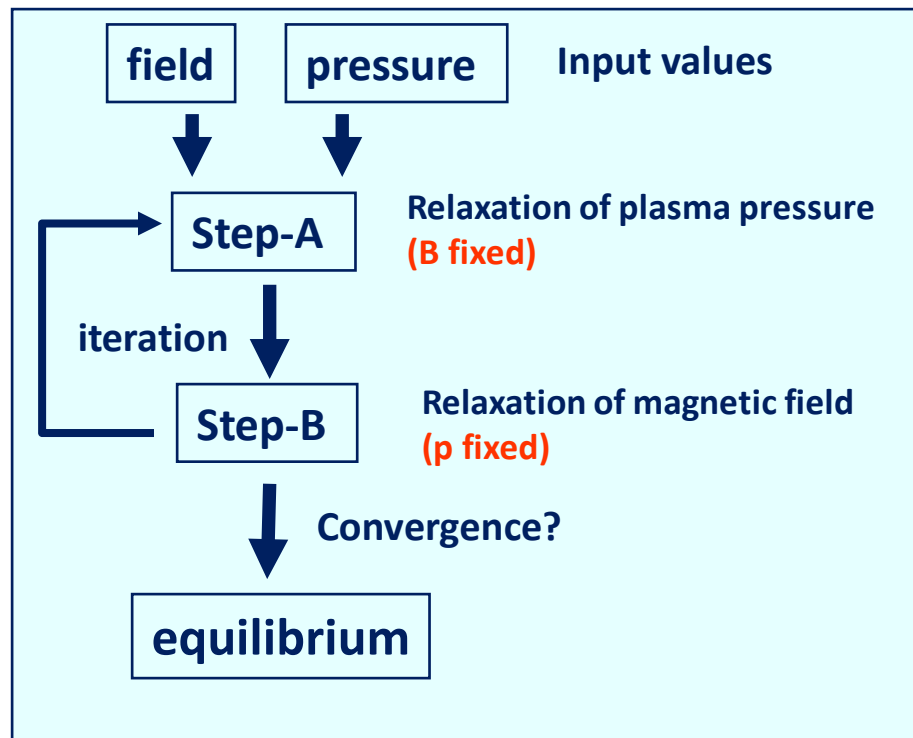
- 1. Nested flux surfaces**
- 2. Stochastic layer (long L_c)**
- 3. Scrape-Off layer (short L_c)**

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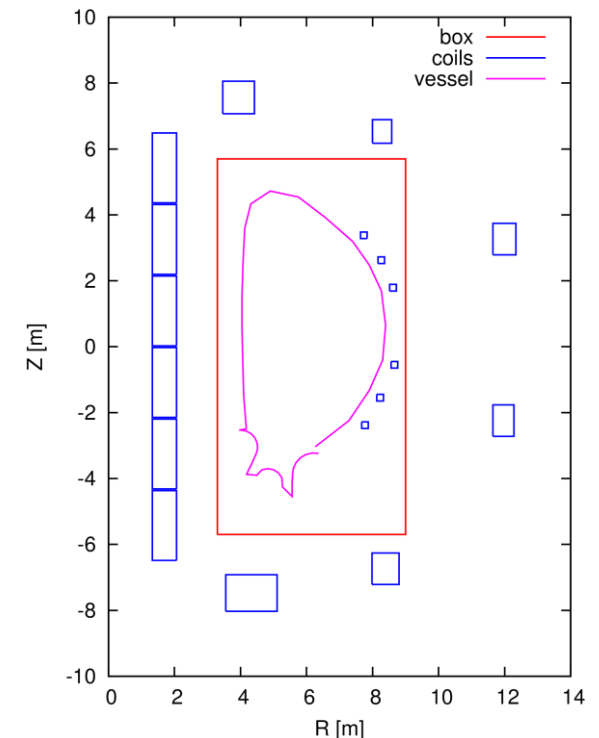
HINT : a 3D MHD equilibrium calculation code

HINT is a 3D MHD equilibrium calculation code without assumption of nested flux surfaces. \Leftrightarrow VMEC (assuming nested flux surfaces)

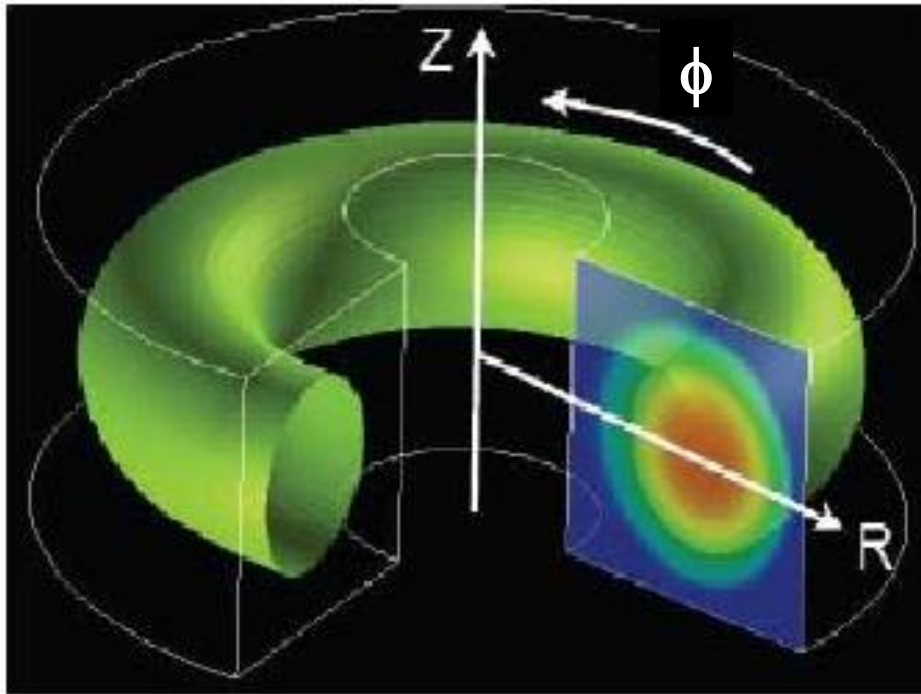
- relaxation method (initial value problem)
- Eulerian coordinate(cylindrical coordinate (R, ϕ, Z))



An example



Numerical model



[Mizuguchi 2012]

- Rectangular Grid on the cylindrical coordinates (R, ϕ, Z)
- 4th order finite difference
- 4th order Runge-Kutta method for time integration

Our approach is the DNS (Direct Numerical Simulation).

Step-A : Relaxation of plasma pressure

This process calculates constant pressure along field line. (**B fixed**)

$$\mathbf{B} \cdot \nabla p = 0 \quad \longleftrightarrow \quad p^{i+1} = \bar{p} = \frac{\int_{-L_{in}}^{L_{in}} \mathcal{F} p^i \frac{dl}{B}}{\int_{-L_{in}}^{L_{in}} \frac{dl}{B}}, \quad \mathcal{F} = \begin{cases} 1 & : \text{ for } L_C \geq L_{in} \\ 0 & : \text{ for } L_C < L_{in} \end{cases}$$

HINT modeling

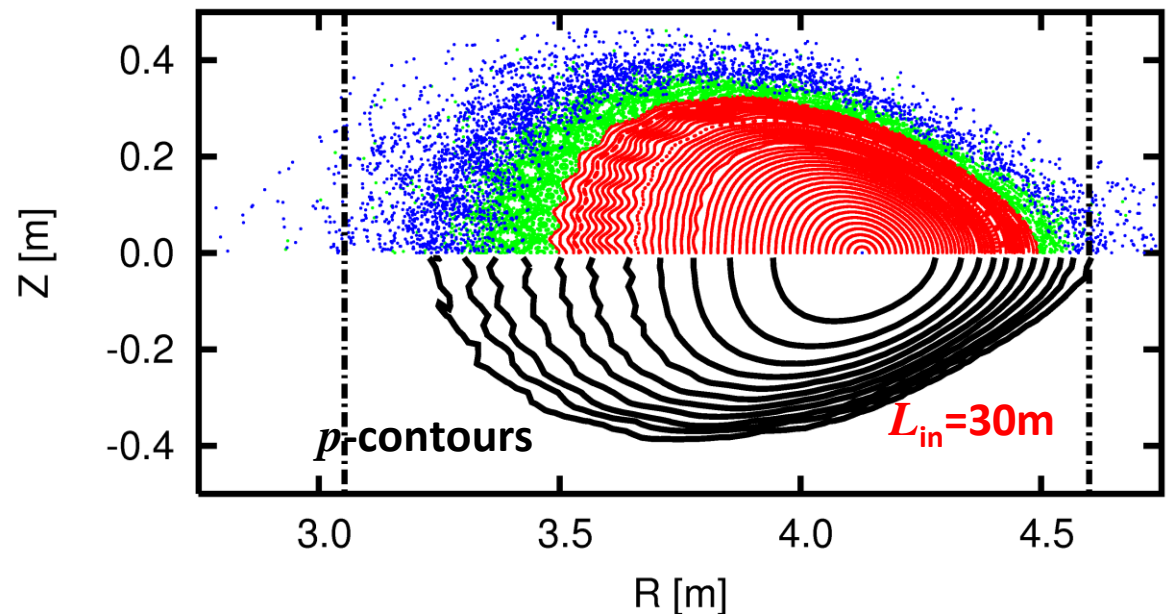
Averaged pressure is calculated by the following equation;

If L_C is longer than λ_e , stochastic field lines keep the pressure.



The structure due to the averaged pressure appears.

$$R_{ax}=3.85, B_T=3.0, B_Q=100\%, \gamma=1.254 \quad \langle \beta \rangle = 2.1\%$$



From the viewpoint of the transport, the electron is sensitive within λ_e because of perpendicular transport by Coulomb collision. Contour lines means averaged structure of field lines.

Step-B : relaxation of magnetic field

This process calculates the time evolution of dissipative MHD equations.

(p fixed)

$$\begin{aligned}\frac{\partial \mathbf{v}_1}{\partial t} &= -\nabla p + \mathbf{j} \times (\mathbf{B}_0 + \mathbf{B}_1) + \mu \Delta \mathbf{v}_1 \\ \frac{\partial \mathbf{B}_1}{\partial t} &= \nabla \times [\mathbf{v}_1 \times (\mathbf{B}_0 + \mathbf{B}_1) - \eta (\mathbf{j}_1 - \mathbf{j}_{\text{net}})] \\ \mathbf{j}_1 &= \nabla \times \mathbf{B}_1\end{aligned}$$

If $d\mathbf{v}_1/dt$ and $d\mathbf{B}_1/dt \Rightarrow \mathbf{0}$, calculation is convergence!

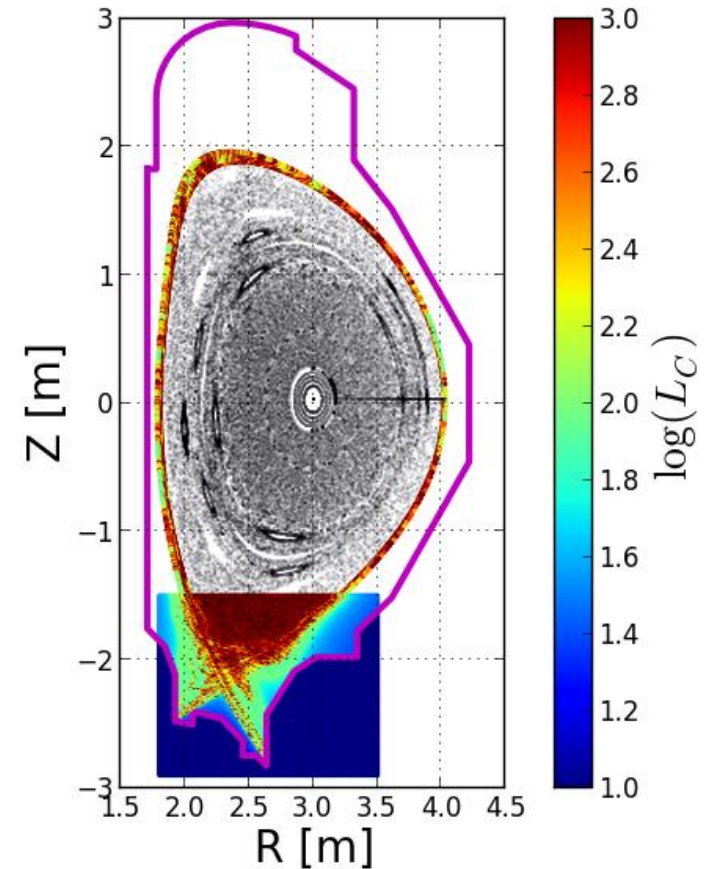
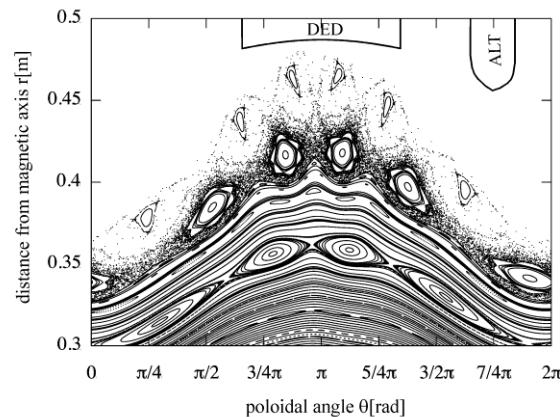
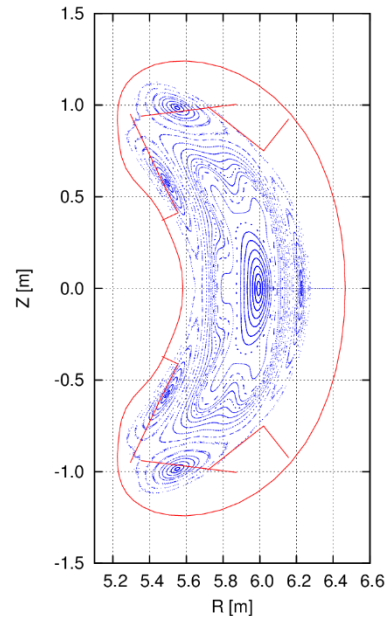
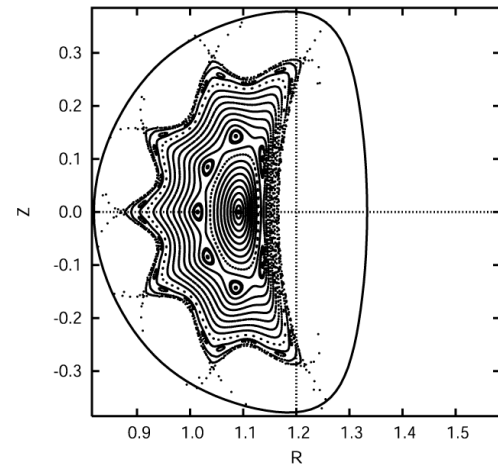
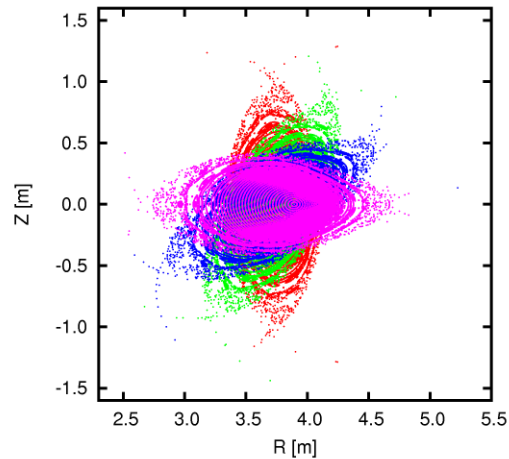
$\rho = 1$ (constant)

$\mathbf{B} = \mathbf{B}_0$ (vacuum) + \mathbf{B}_1 (equilibrium response)

$\mathbf{v} = \mathbf{v}_1$ (MHD flow)

η : resistivity (constant)

HINT applies to many stellarator and tokamaks



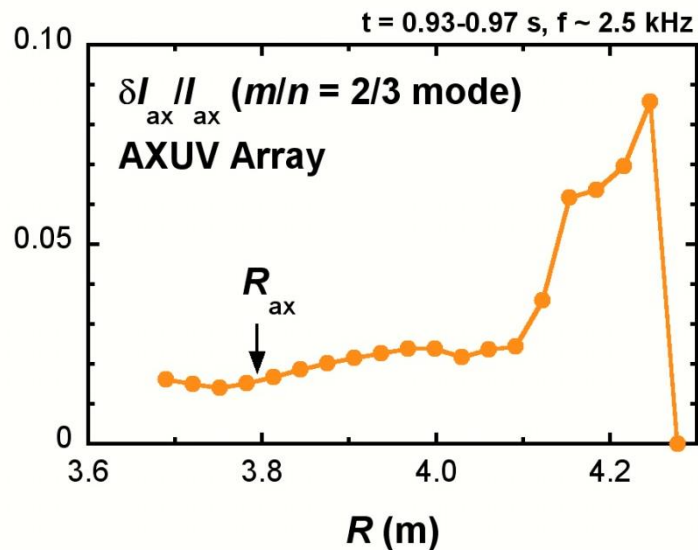
TJ-II, ITER, DIII-D, JET, MAST,...

EAST, ASDEX-U, (planning)

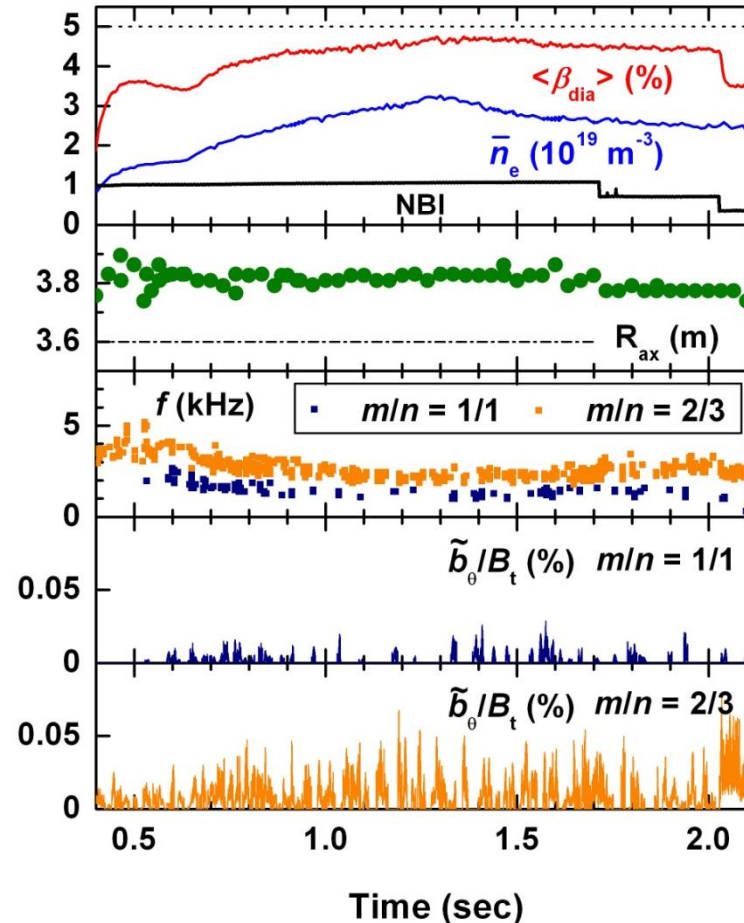
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High-beta Steady State Discharge

- ▶ $\langle \beta \rangle_{\max} \sim 4.8\%$, $\beta_0 \sim 9.6\%$, $H_{\text{ISS95}} \sim 1.1$
- ▶ Plasma was maintained for $85\tau_E$
- ▶ Shafranov shift $\Delta/a_{\text{eff}} \sim 0.25$
- ▶ Peripheral MHD modes are dominantly observed.

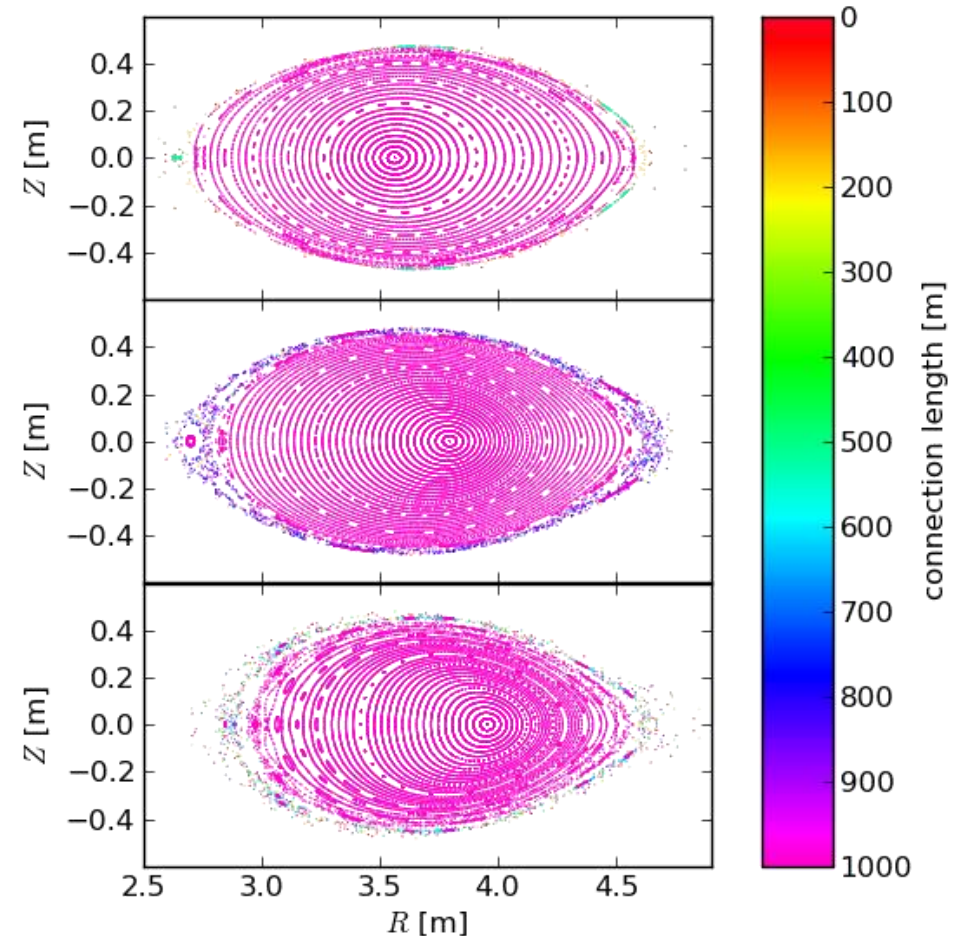
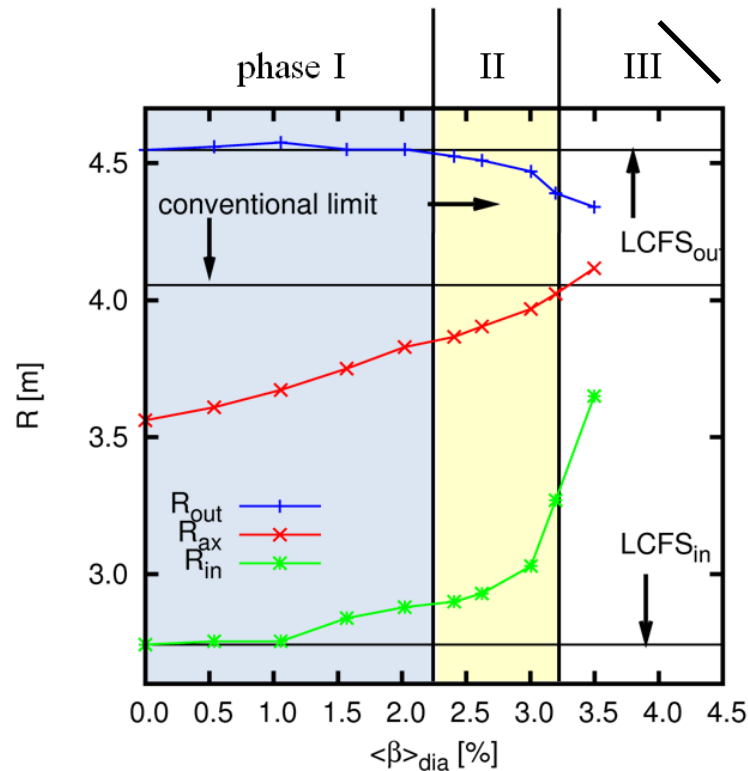


$$R_{\text{ax}} = 3.6\text{ m}, B_t = -0.425\text{ T}$$



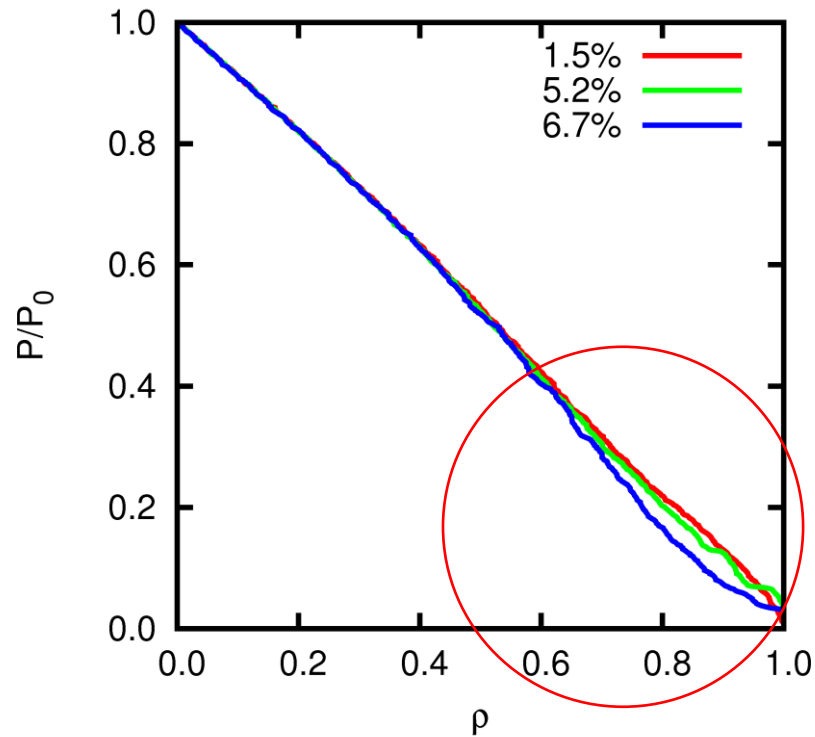
How about is the magnetic surface topology?

Degradation of flux surfaces due to increasing β

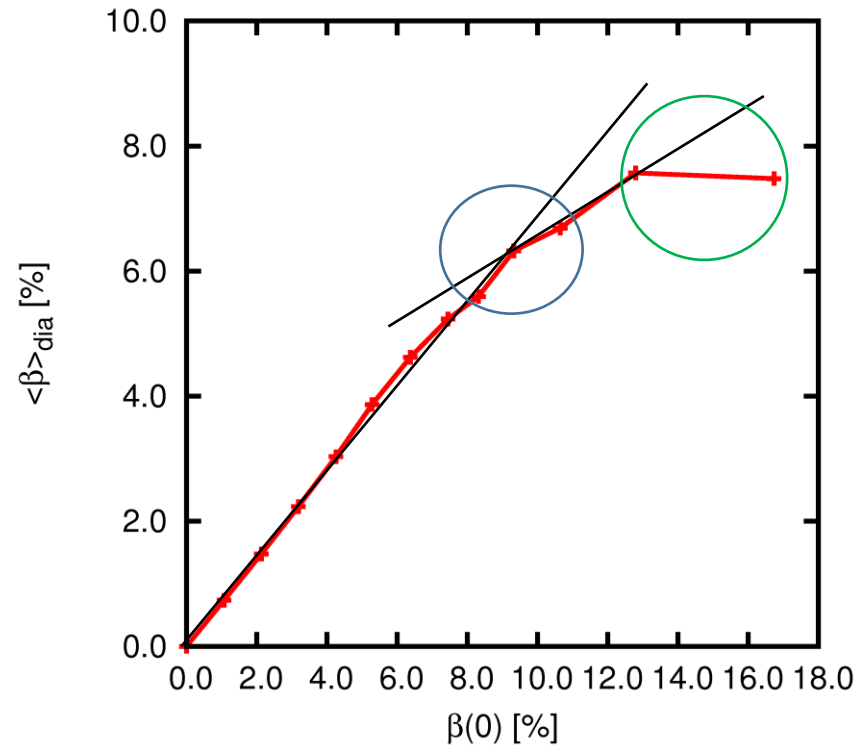


In the peripheral region, magnetic field lines become stochastic as β increases. The volume inside LCFS shrinks drastically.

Results of HINT analyses



For $\beta > 6.7\%$, the fixed pressure profile is reduced at $\rho > 0.6$.



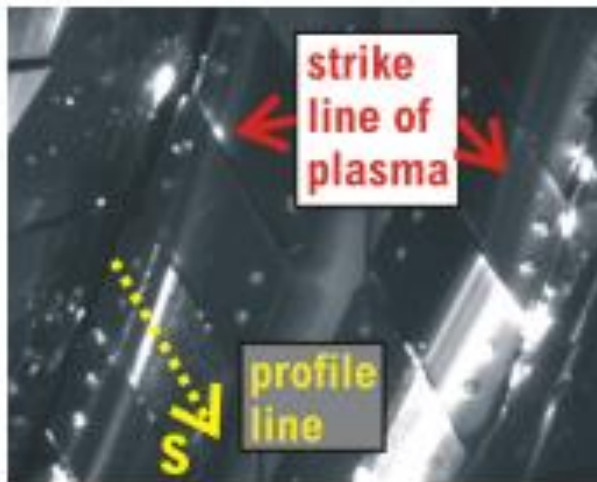
The slope of $\langle \beta \rangle$ changes due to the reduction of the pressure profile.

**The change of the slope is a good index of the equilibrium beta limit.
Note: this index is a soft limit of the MHD equilibrium.**

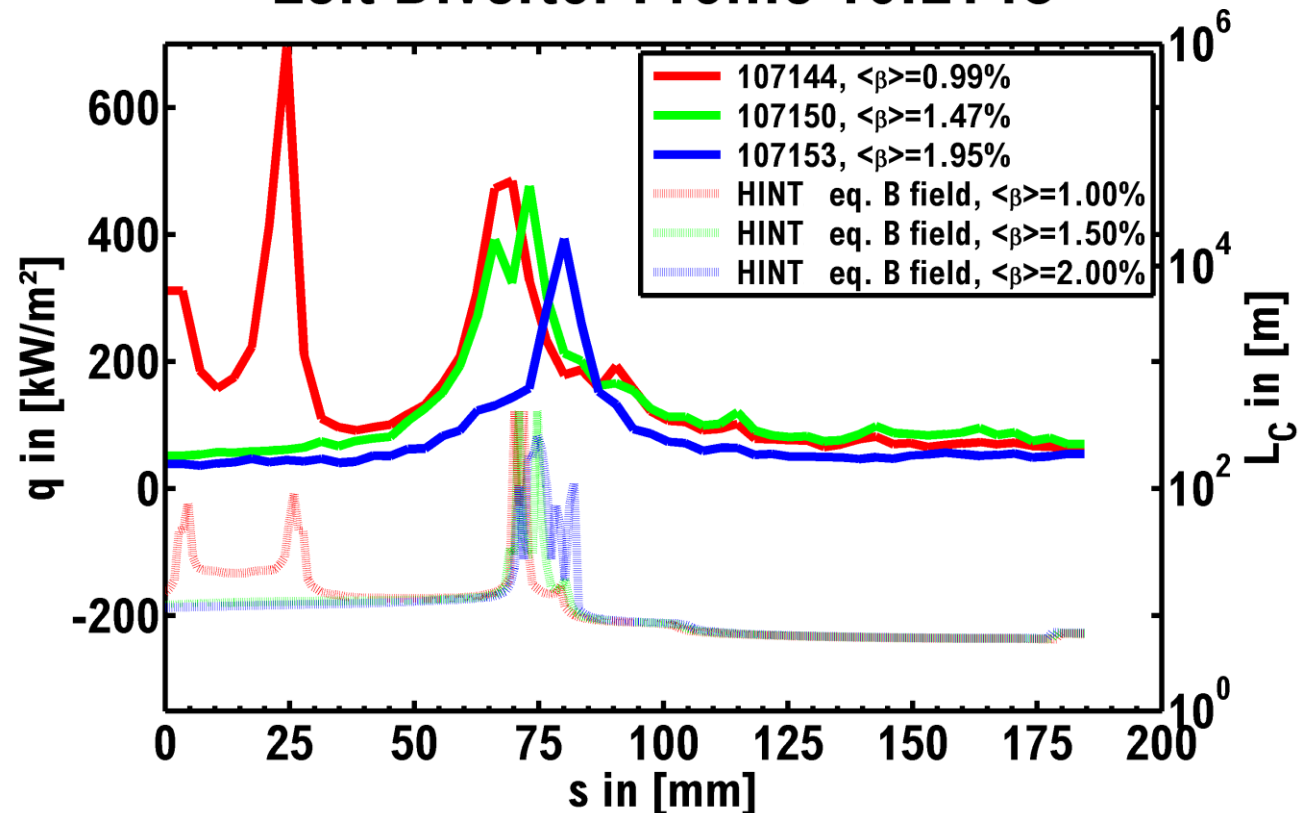
Heat flux strongly depends on magnetic topology

Stochastization of magnetic field lines strongly affects on divertor heat load.

In LHD, divertor heat load was measured by IR camera system.



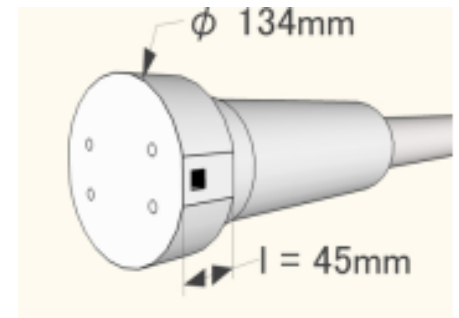
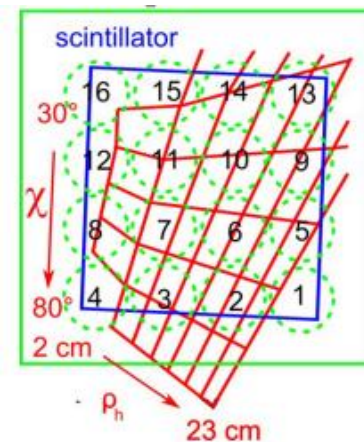
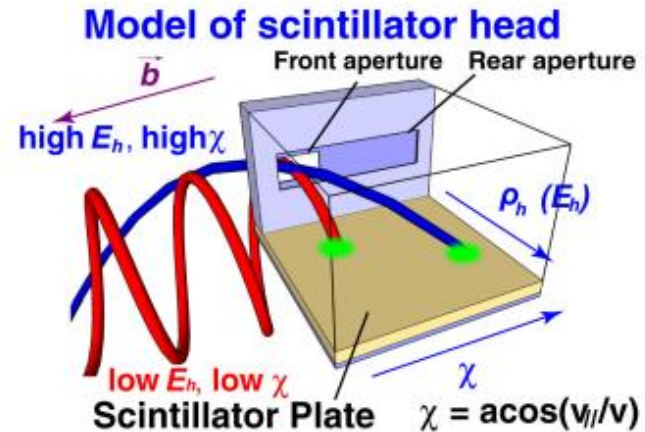
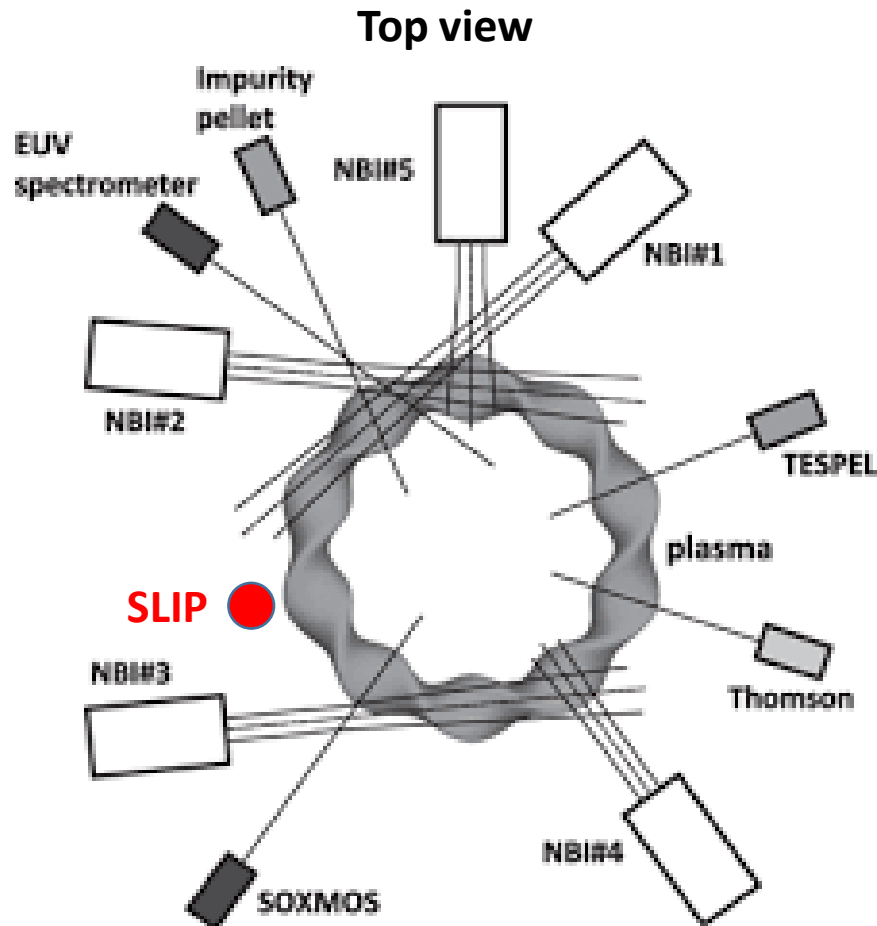
Left Divertor Profile 10IL14C



HINT can give reliable prediction!

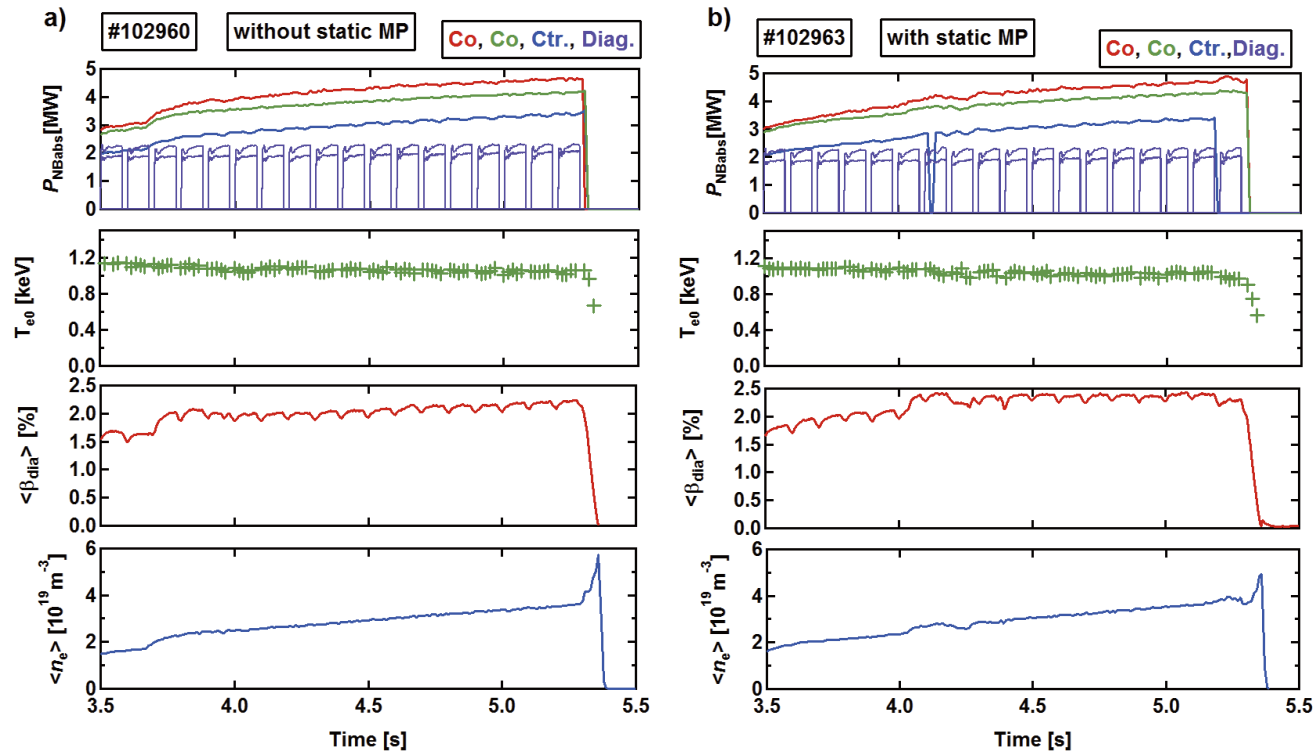
Fast ion orbit is very sensitive to magnetic topology

Scintillator-based Lost Ion Probe (SLIP) is installed to measure fast ion loss.



RMP experiment in LHD

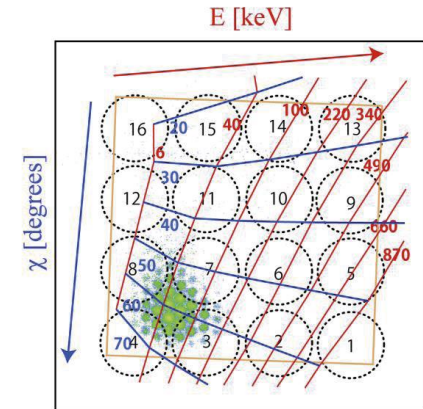
Standard configuration (Rax=3.6m), B=-0.9T



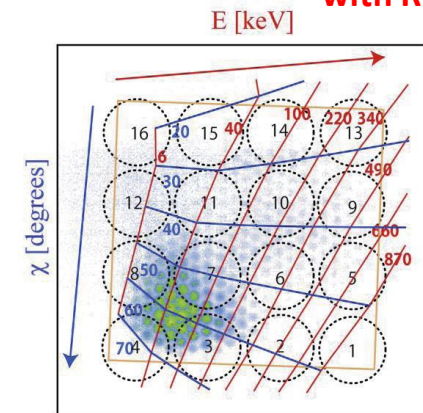
Almost same shots could be reproduced with and w/o RMP.

Scintillator image (SLIP, 8-O)

a) #102960 t=4.50 s **w/o RMP**



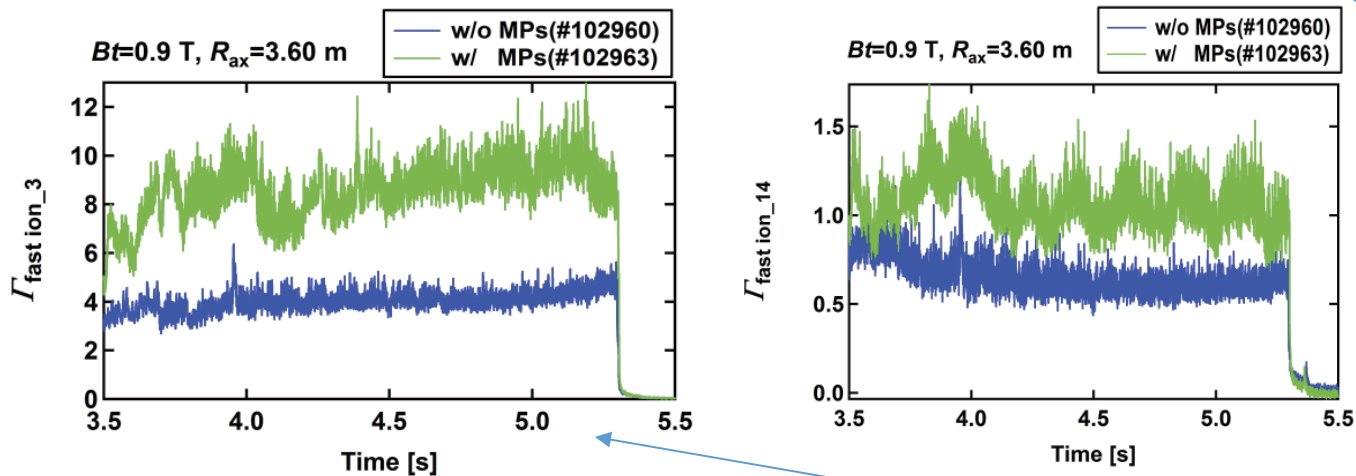
b) #102963 t=4.50 s **with RMP**



RMP experiment in LHD

Time evolutions of fast-ion-loss signals measured by SLIP at PMT 3 and 14.

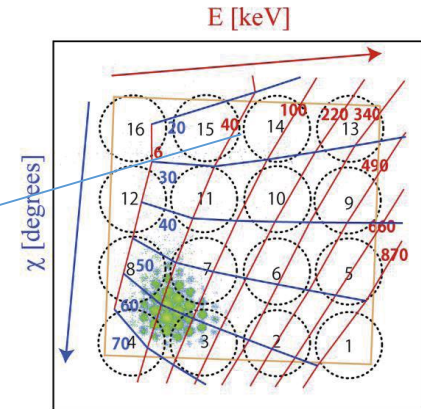
RMP enhances fast ion loss.



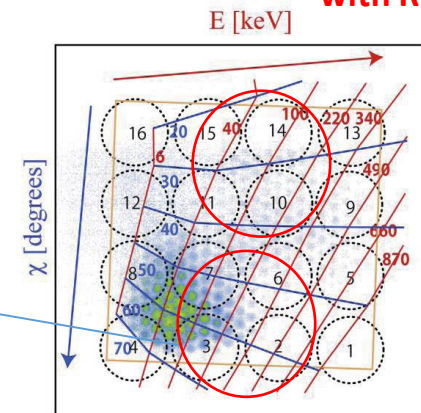
- Low-pitch angle/high energy loss increases.
- High-pitch angle/low energy loss increases.

Scintillator image (SLIP, 8-O)

a) #102960 $t=4.50 \text{ s}$ **w/o RMP**

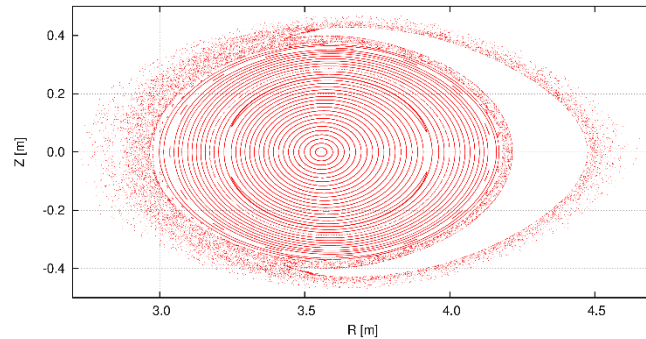


b) #102963 $t=4.50 \text{ s}$ **with RMP**

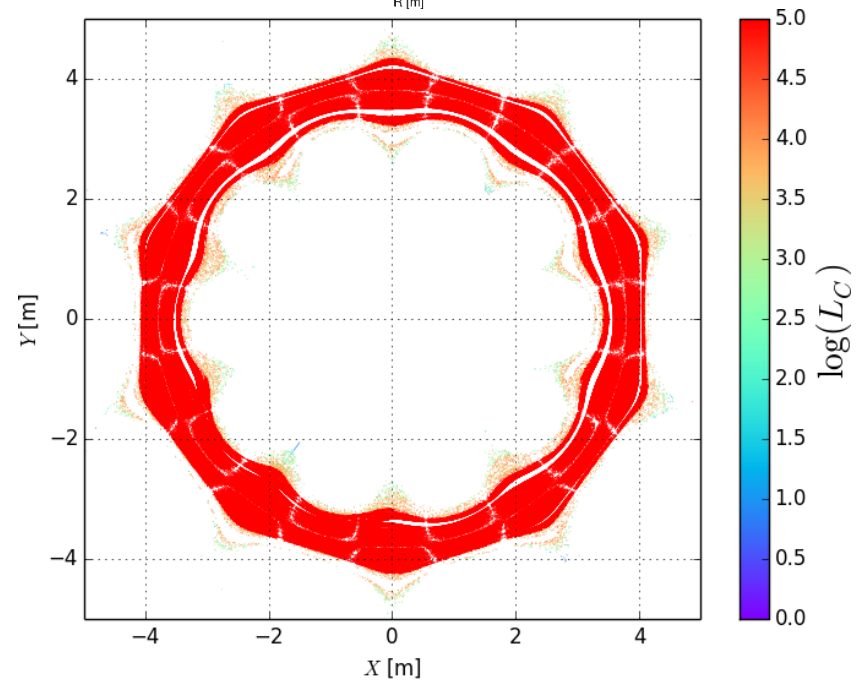
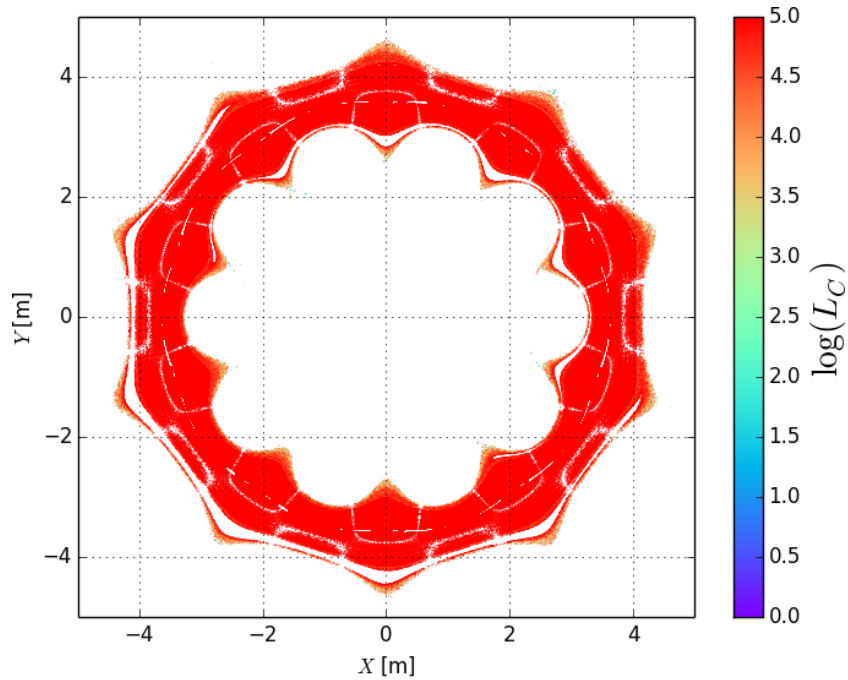
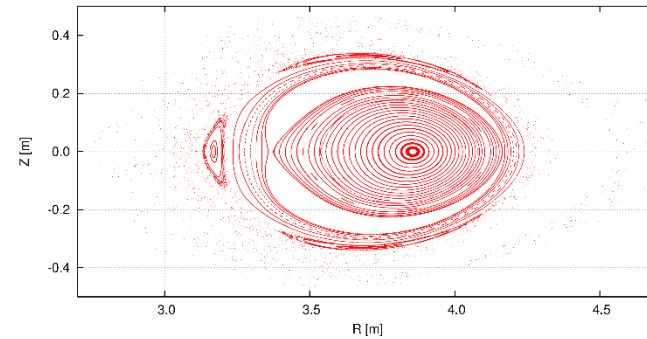


Impact of beta effects on magnetic topology

vacuum

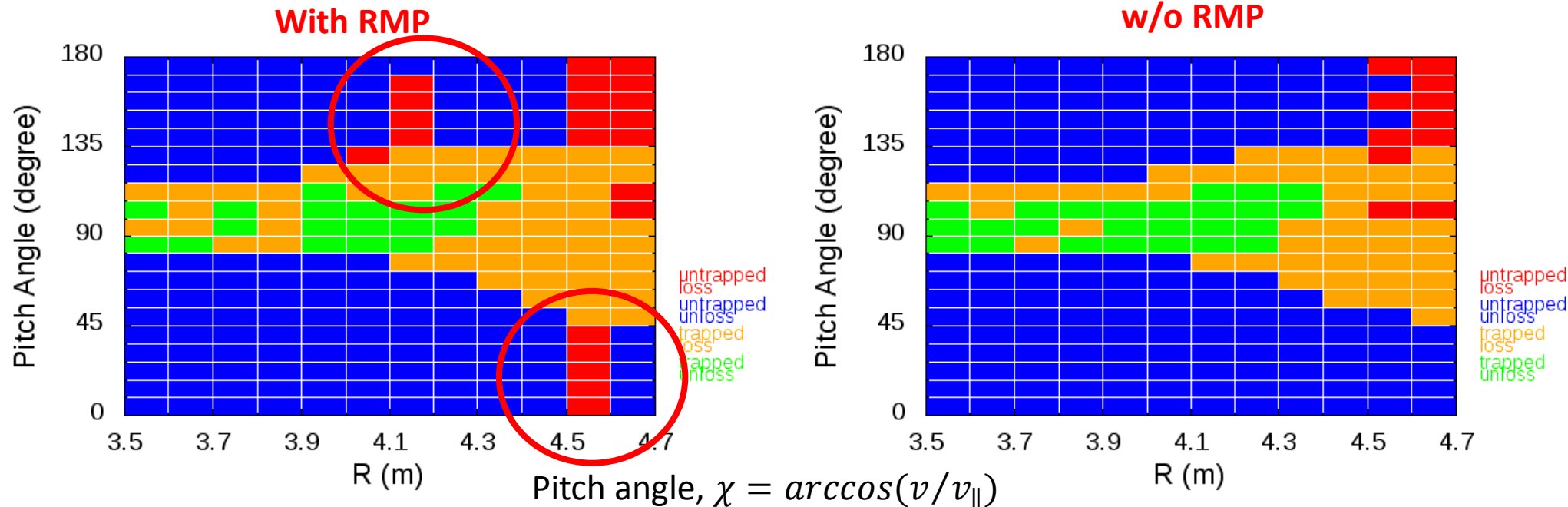


$\langle \beta \rangle \sim 2.4\%$



Impacts of RMP on fast ion confinement

To study impacts of RMP qualitatively, collisionless orbit of fast is studied.
 $E = 90\text{keV}$, ϕ and Z of starting point are fixed to 234 deg and 0 m, respectively.



Because of RMP...,

1. Loss of passing particles starting from $R \sim 4.5$ m (Co) and $R \sim 4.1$ m (Cntr) is increased.
2. Loss of trapped particle starting from $R \sim 4.1$ m is increased.

These are consistent to SLIP observation.

Courtesy to J. Morimoto

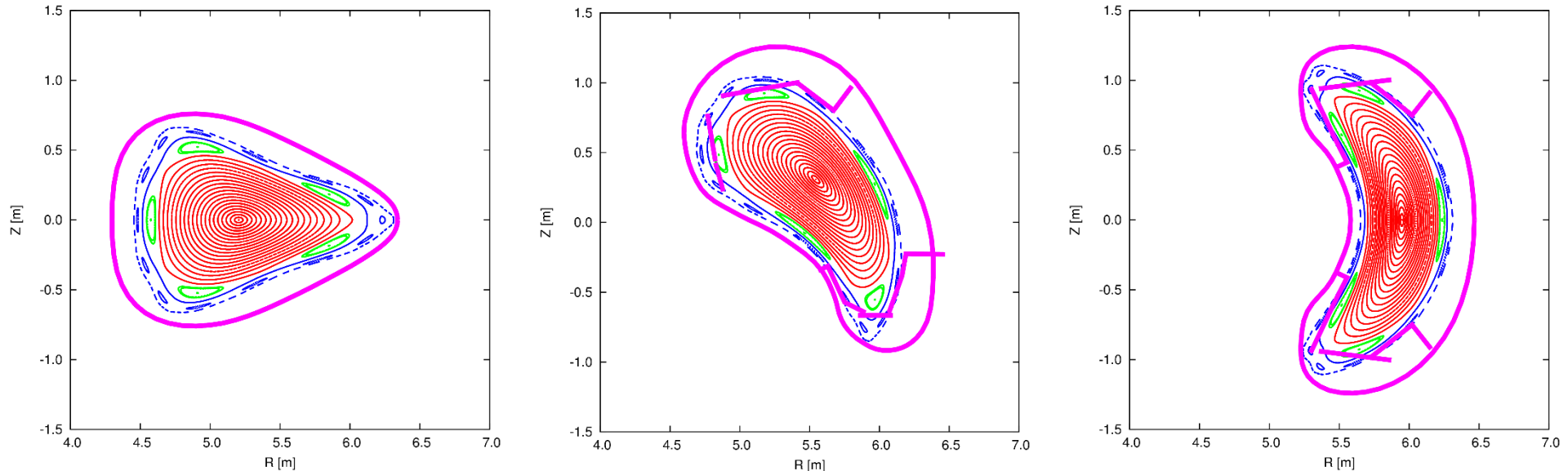




Wendelstein 7-X

- Island divertor concept -

Island divertor: small pitch allows use of perpendicular transport of magnetic islands.



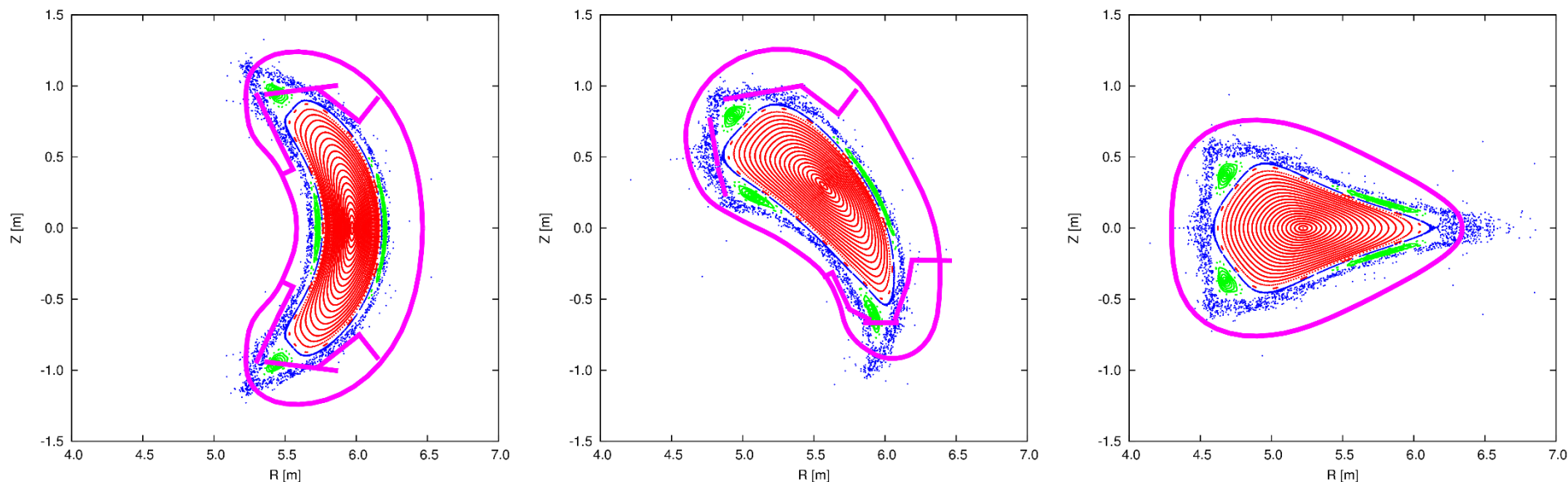
- Optimization of W7-X: MHD-equilibrium, stability, ... **but not for divertor!**
- Natural islands used for separatrix generation to divert fluxes to appropriately designed divertor plates!
- Low shear: Islands generated with small field components ($B_r/B \sim 10^{-4}$)!



Divertor configurations are very sensitive plasma generated fields!
(e.g. beta-effects and net-toroidal current densities)

High-iota configuration - magnetic field configuration -

Divertor islands at iota = 5/4



Strong stochastization is expected.

1. E. Strumberger, et al., in Nuclear Fusion
2. Y. Suzuki, et al., in EPS2006

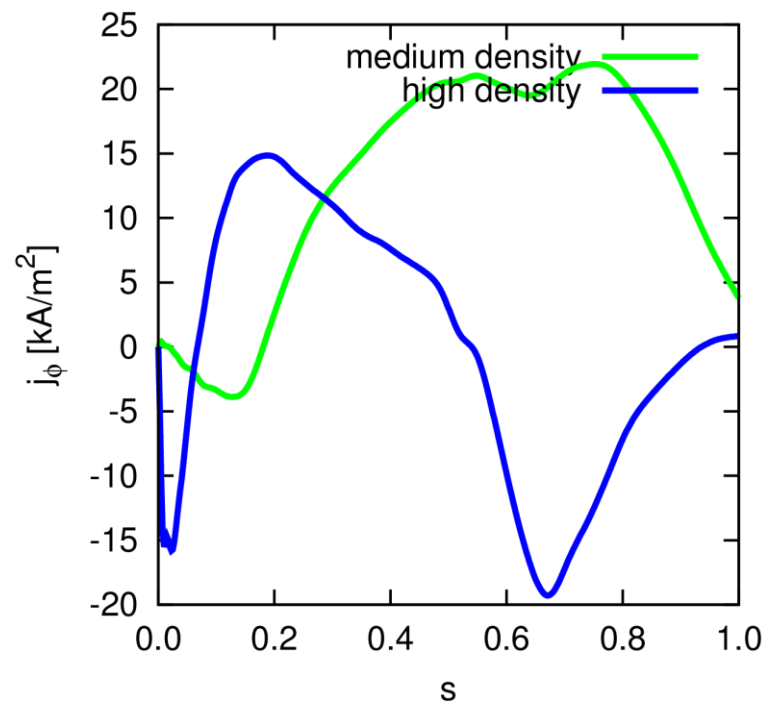
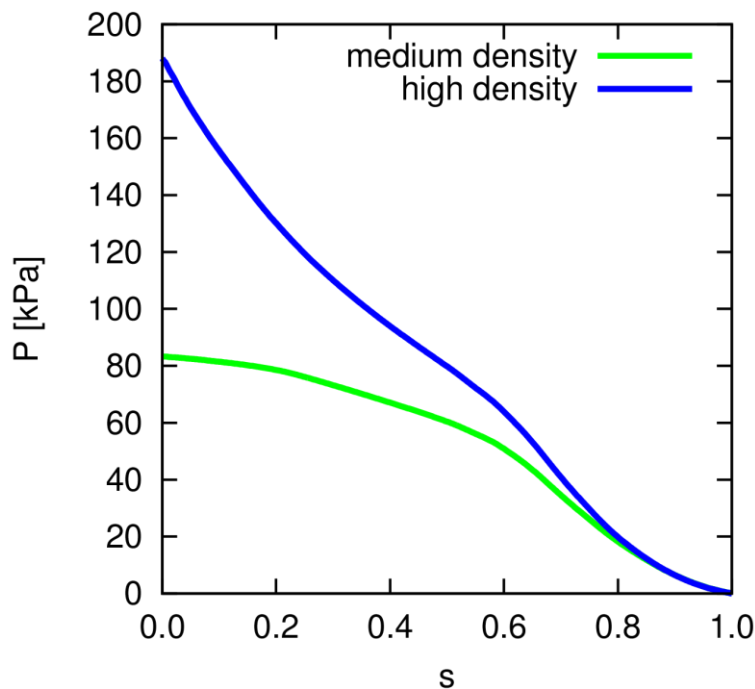
But, the toroidal current density effect not yet studied.



High-iota configuration

- pressure and current density profiles -

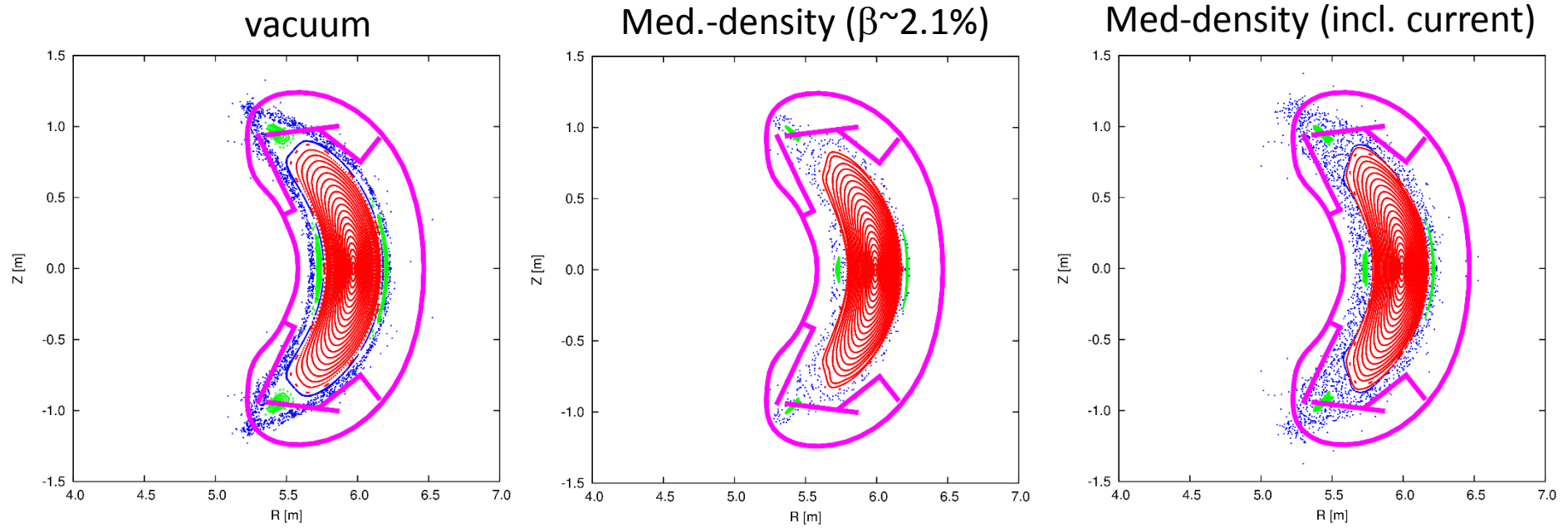
Radial profiles predicted by the self-consistent transport analysis



- Low to high density => broad to more peaked (<=> off- to on-axis heating).
- Electron-root for low and medium density => enhanced bootstrap current density
- High density => small bootstrap current densities ($I_p < 1\text{kA}$).



Impacts of β -sequence and current density - magnetic field configuration -



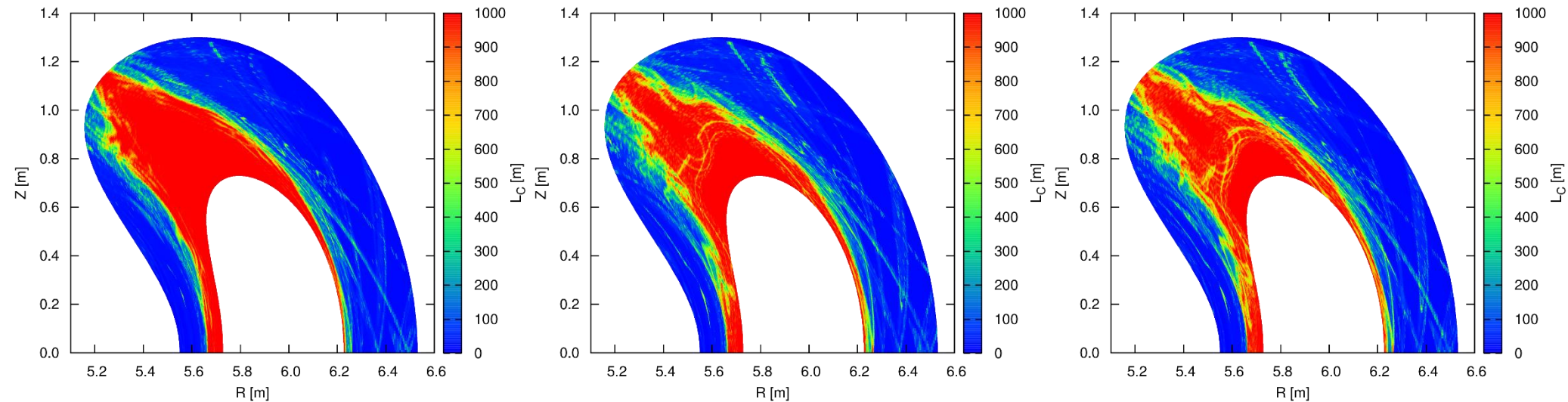
- **5/4 islands are robust up to 3%. But, increasing stochastization.**
- **Higher-order rational islands grow and overlap (e.g. 15/13) \Rightarrow volume shrinks.**
- **Toroidal current densities can shrink island sizes of high-order rationals but enhance stochastization.**
- **For high-density case, 5/5 islands appear in plasma core(beta + bootstrap current density).**
- **higher elongation of flux surfaces \Rightarrow less tight fitting divertor**

Impacts of β -sequence and current density - magnetic field configuration -

vacuum

Med.-density ($\beta \sim 2.1\%$)

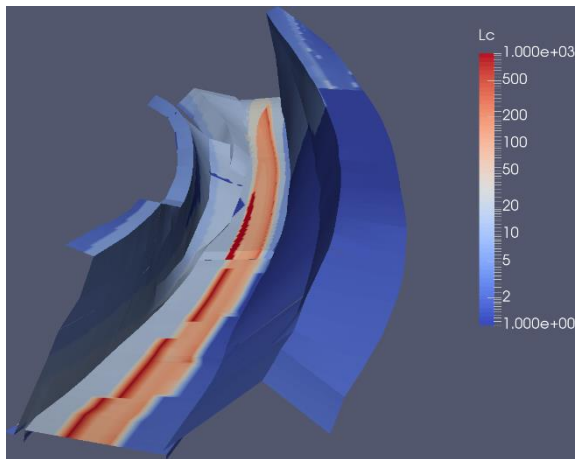
Med-density (incl. current)



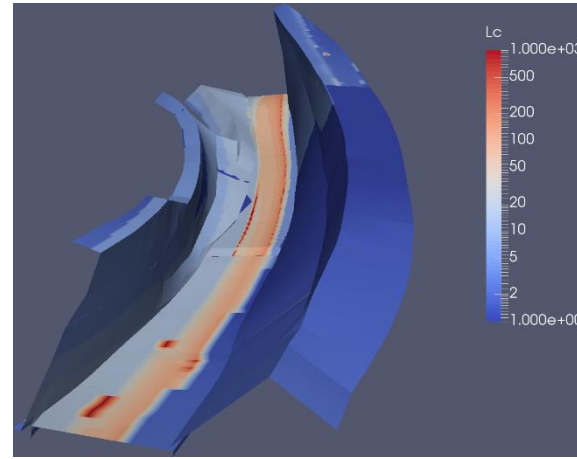
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Impacts to divertor operation - footprints onto in-vessel components --

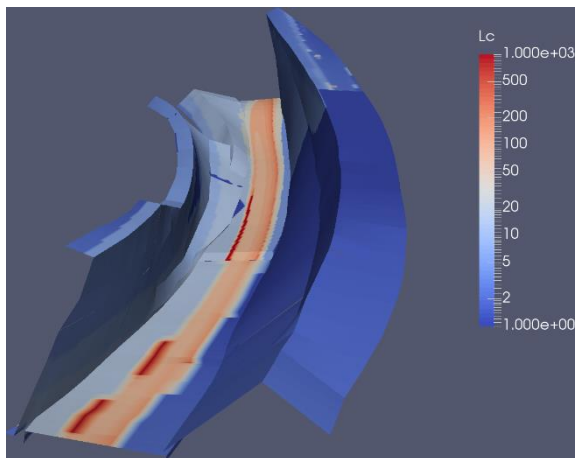
vacuum



Medium density



Medium density incl. current

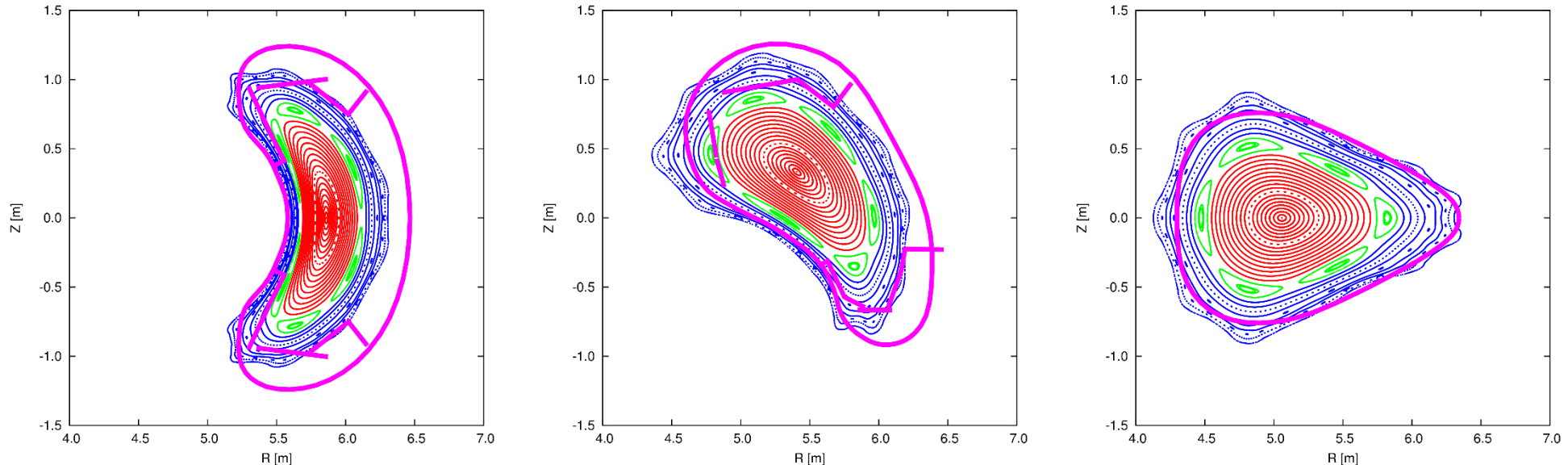


- For vacuum case, plasma is diverted onto in-vessel components.
- For medium-density case, connection length becomes short because of enhanced stochasticity.
- For including current density, plasma is still diverted by in-vessels and new splitting appear.

Minimum bootstrap current configuration at low-iota

Minimum bootstrap current at low-iota \Leftrightarrow large toroidal mirror field $b_{01} \approx 24\%$

Divertor islands at iota = 5/6



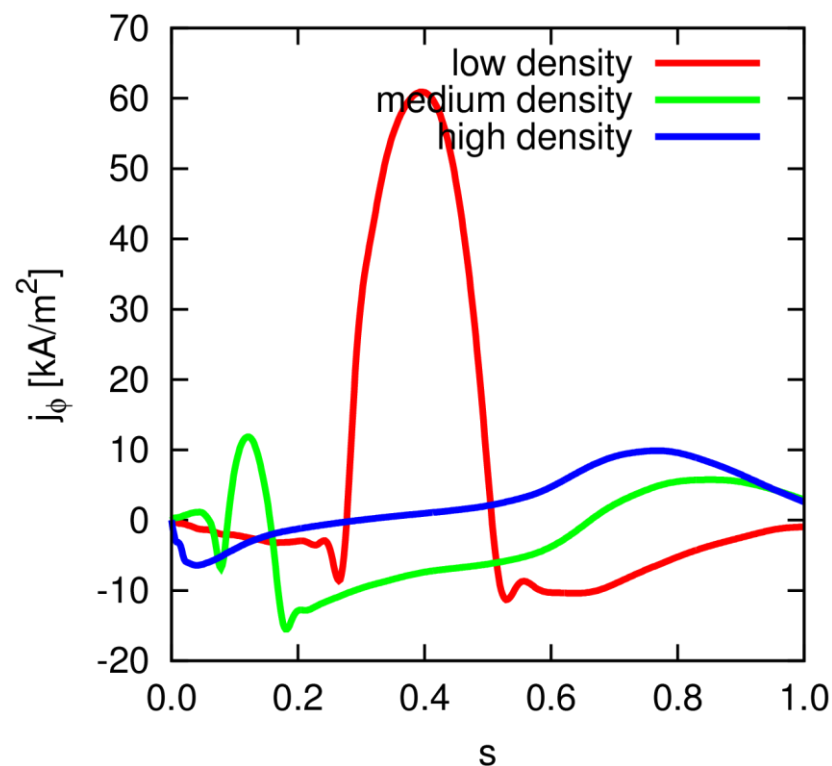
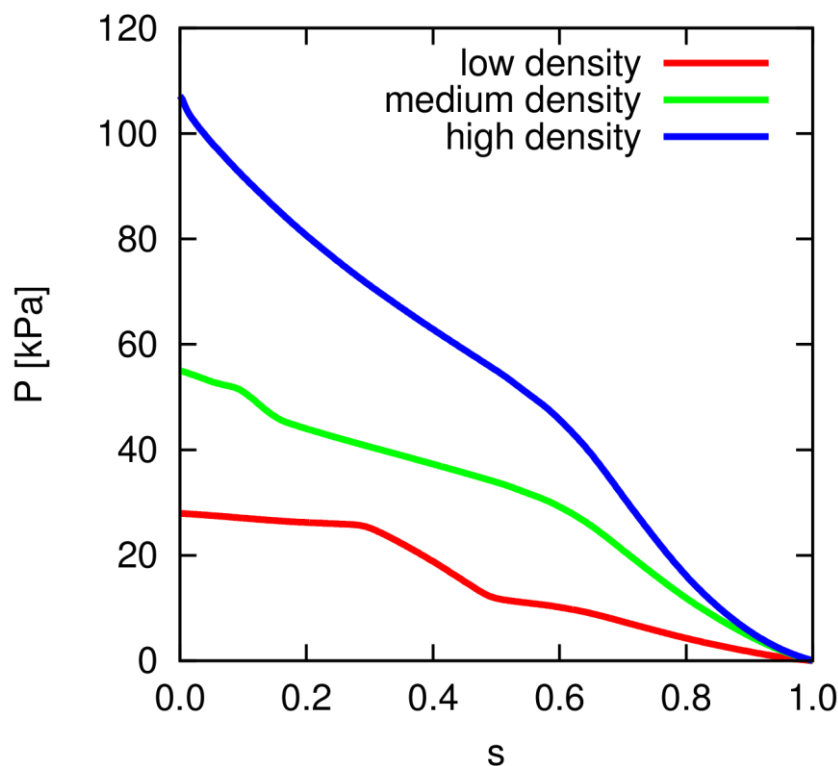
- Strong toroidal variation of cross section due to large mirror field
small cross section in bean-shaped plane
divertor interaction around $\phi=16^\circ$ (teardrop cross section)
divertor geometry adapted for mirror fields from 0% to 10%
- Self-consistent transport simulations focussed on core confinement (J. Geiger, et al., EPS2014). Adjustments of configuration might be needed for divertor compatibility.



Low-iota configuration

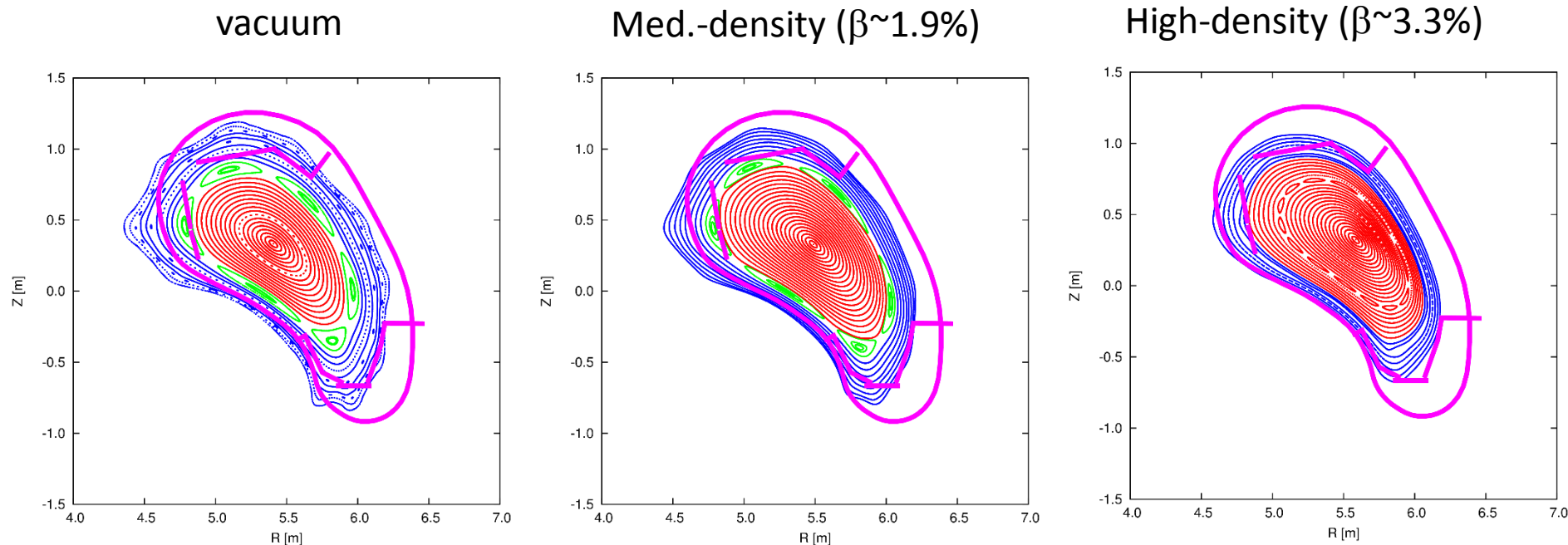
- pressure and current density profiles -

Radial profiles predicted by the self-consistent transport analysis



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Impacts in β -sequence - magnetic field configuration -

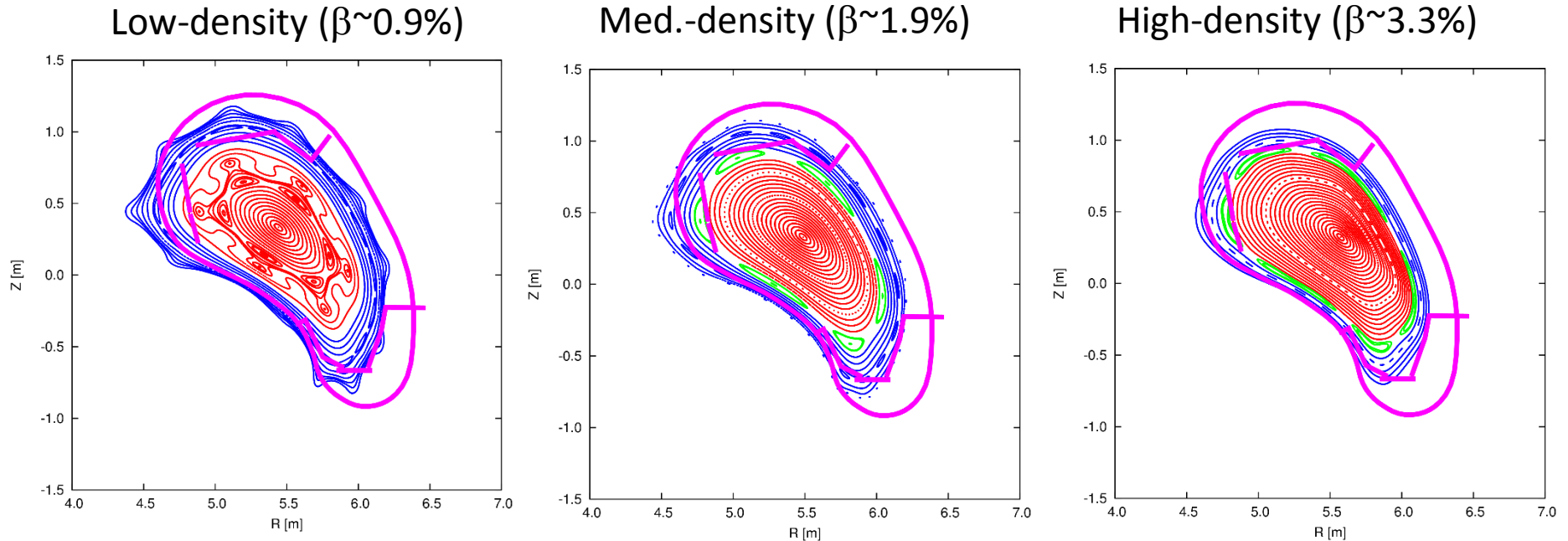


- 5/6 islands are robust for vacuum and low density cases.
- Medium density case \Rightarrow 5/6 Islands shift radially outward \Rightarrow Limiter configuration.
- High density case \Rightarrow 5/6 islands vanish (self-healing), 10/12 islands become visible.
- Plasma is limited by vertical baffle plate in all cases.
- Flux surfaces fill divertor volume rather tight.

\Rightarrow configuration adjustment with planar coils: horizontal positioning, iota-fine-tuning?

\Rightarrow control 5/6-island width by sweep coils?

Impacts of toroidal current density - magnetic field configuration -



Current density profile from transport simulation

5/6 islands are robust.

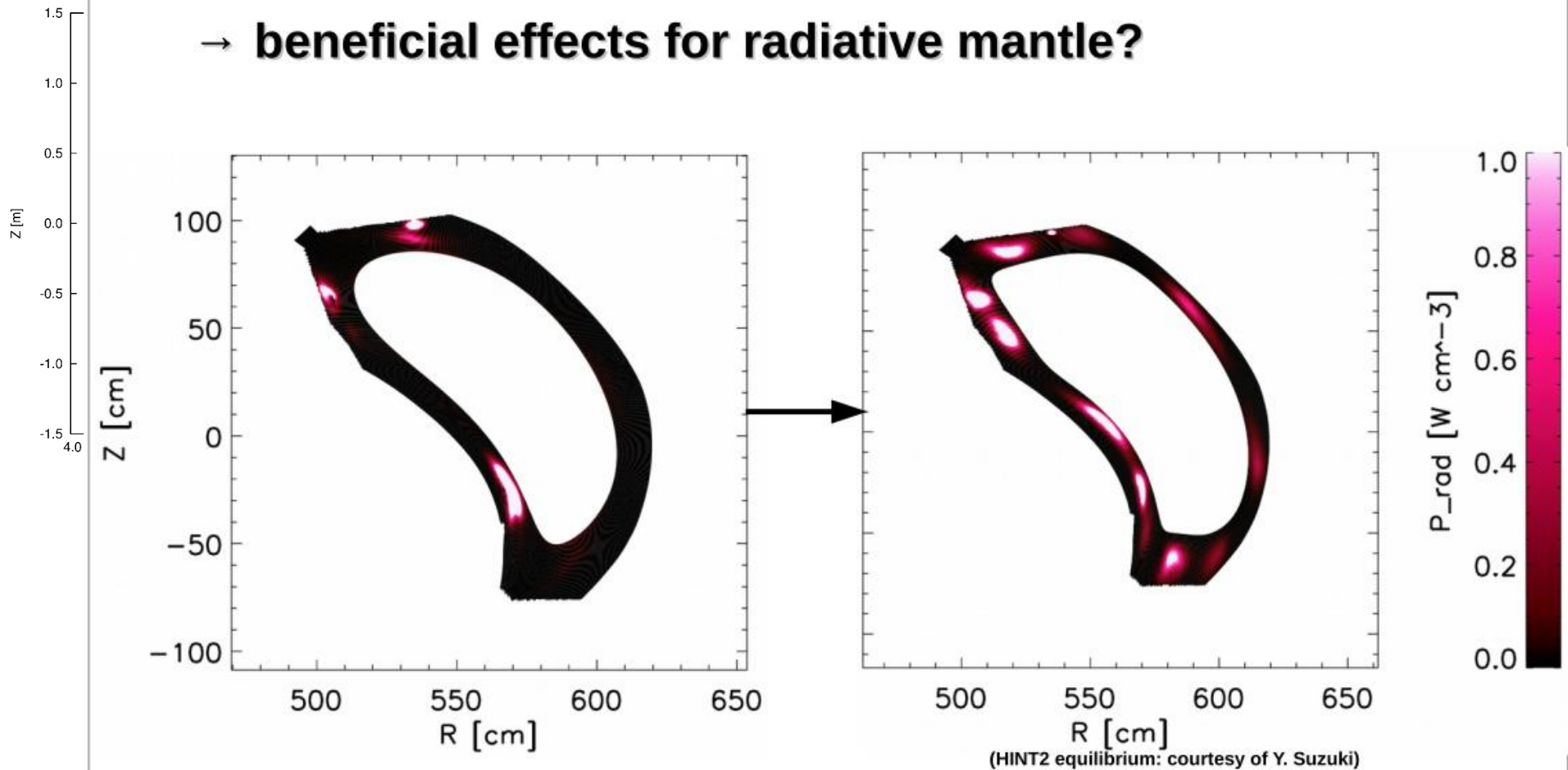
Radial location of 5/6 islands moves according to current densities

- low density: islands inside due to localized bootstrap current density
- medium + high density: islands at boundary, but smaller for high-density

Positive current => diverted, negative current => limited

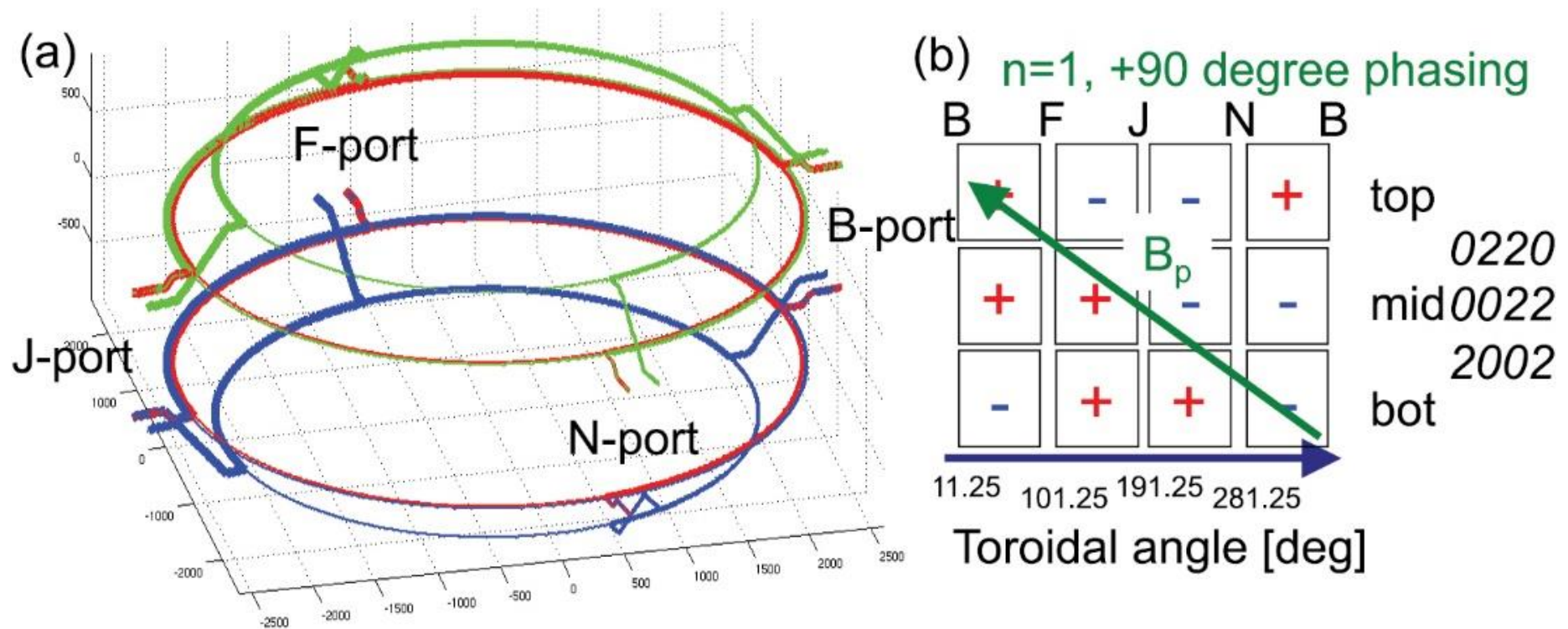
Equilibrium effects seem to cause more homogenous radiation distribution

→ beneficial effects for radiative mantle?

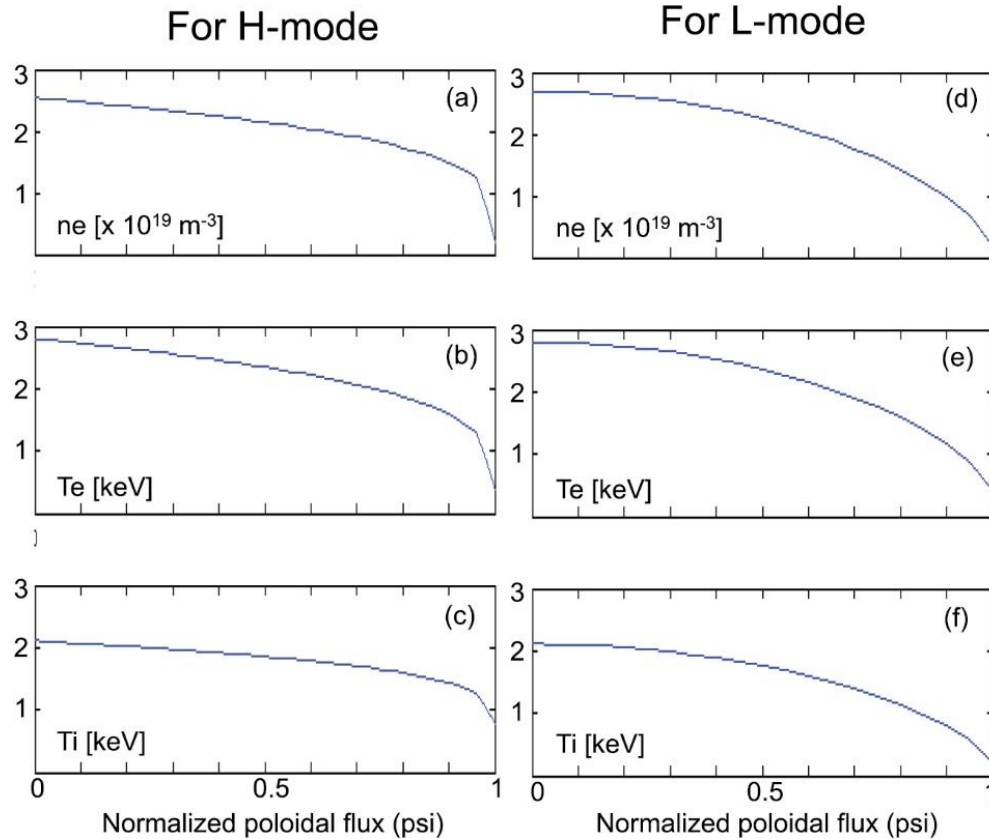
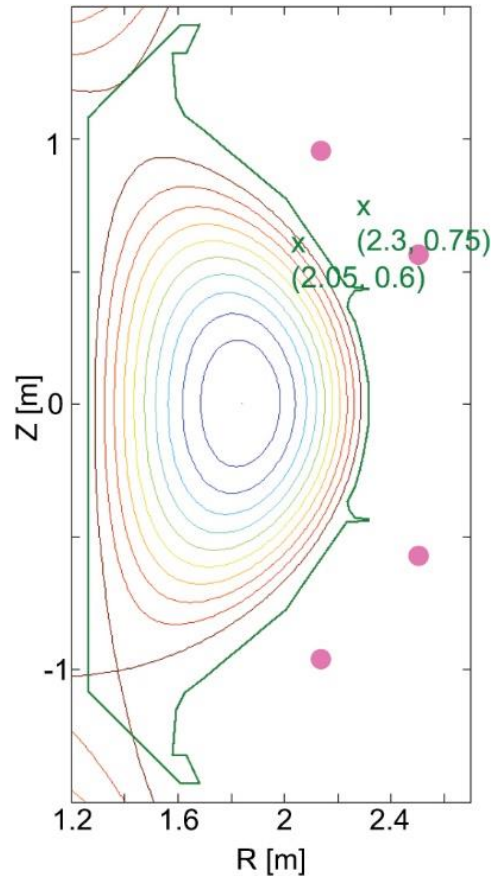


Impact of RMP on magnetic topology in KSTAR

RMP coils model



Nonlinear response studied for 2 scenarios (H- and L-modes)



- ITER-like LSN configuration is studied.
- Plasma profiles are modeled artificially.

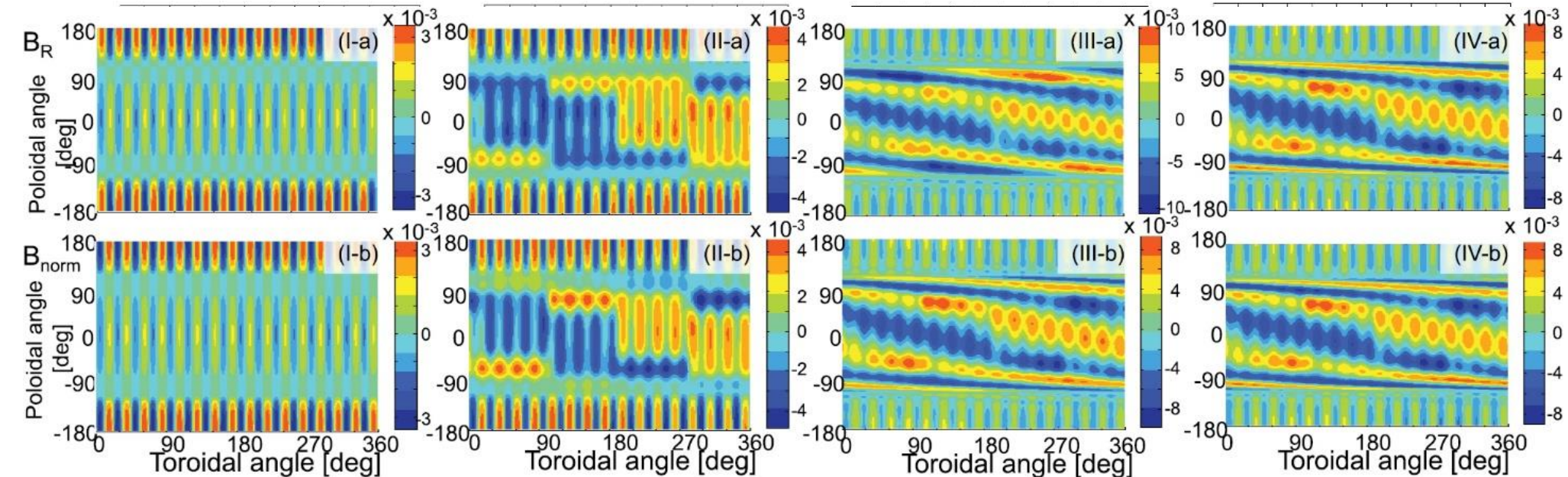
Magnetic field variation in different model and profiles

TF ripple only

Vacuum model

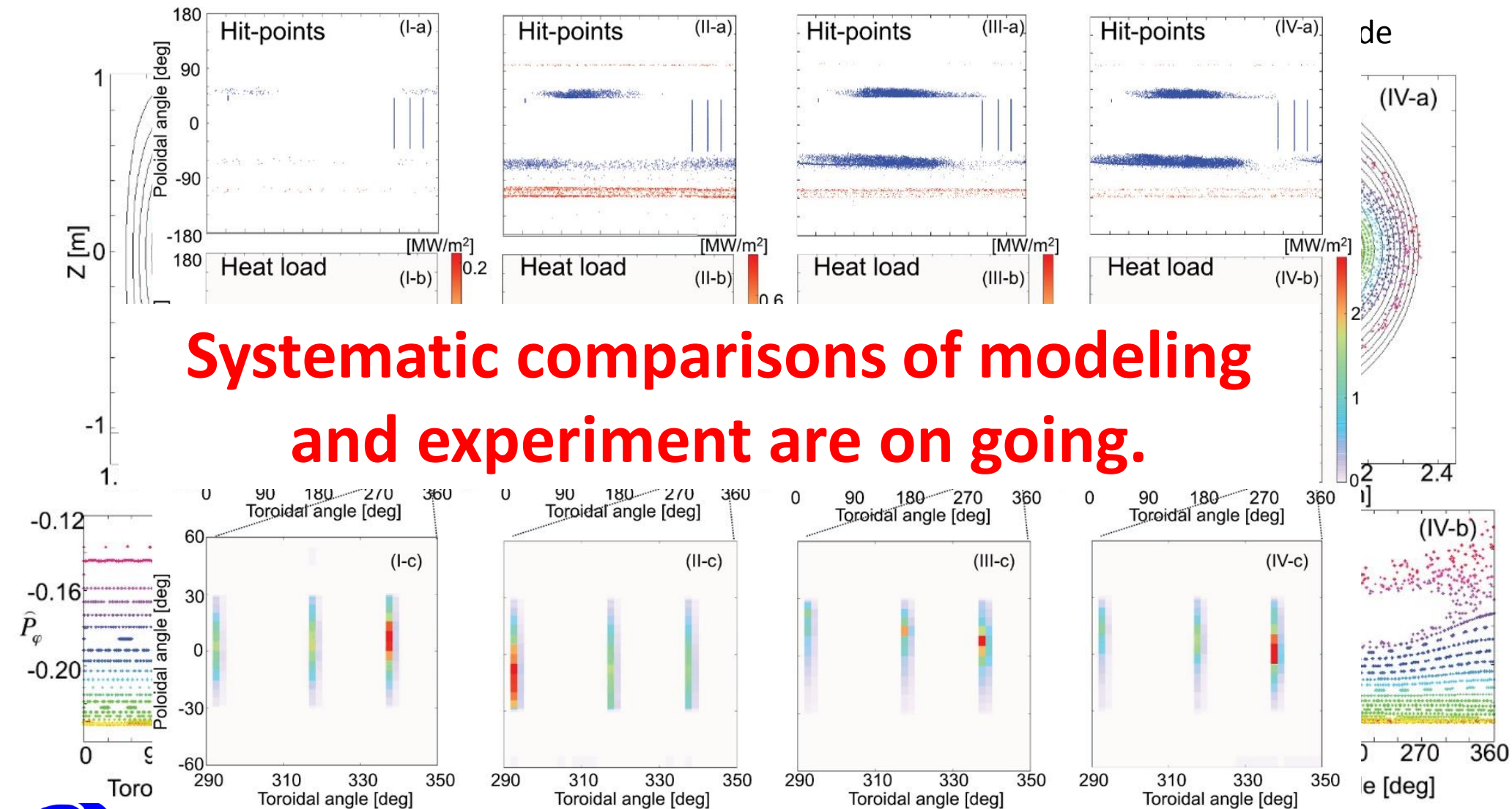
HINT (H-mode)

HINT (L-mode)



- Step-like $n=1$ pattern appeared for vacuum approximation.
- Smooth $n=1$ pattern appeared for HINT, because $n=1$ RMP makes helical distortion of plasma column.

Topological change seems small for L- and H-mode cases, but...



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Impact of plasma rotation on 3D MHD equilibrium

In many experiments, the plasma rotation is observed.



To simulate more realistic 3D equilibrium, the rotation should be included. And, the impact of the rotation to the magnetic field is interesting.

- **Assuming only toroidal rotation.**
- **Assuming isothermal condition.**
- **If the toroidal plasma rotation exists, the pressure distribution is shifted by the inertial force.**

$$p \equiv p_0(\Phi) \exp \left[M^2 \left(\frac{R^2}{R_0^2} - 1 \right) \right] \quad \text{Here,} \quad M^2(\Phi) \equiv \left(\frac{v_\phi}{v_{th}} \right)^2 = \frac{m_i R_0^2 \Omega^2}{2T}$$

Improvement of HINT to include plasma rotation

$$\begin{aligned}\frac{\partial \mathbf{v}_1}{\partial t} &= -\mathbf{v}_0 \cdot \nabla \mathbf{v}_0 - \nabla p + \mathbf{j} \times (\mathbf{B}_0 + \mathbf{B}_1) + \mu \Delta (\mathbf{v}_0 + \mathbf{v}_1) \\ \frac{\partial \mathbf{B}_1}{\partial t} &= \nabla \times [(\mathbf{v}_0 + \mathbf{v}_1) \times (\mathbf{B}_0 + \mathbf{B}_1) - \eta (\mathbf{j}_1 - \mathbf{j}_{\text{net}})] \\ \mathbf{j}_1 &= \nabla \times \mathbf{B}_1\end{aligned}$$

Toroidal flow, \mathbf{v}_0 , is prescribed by the function of the toroidal flux.

$\rho = 1$ (constant)

$\mathbf{B} = \mathbf{B}_0$ (vacuum) + \mathbf{B}_1 (equilibrium response)

$\mathbf{v} = \mathbf{v}_0$ (given toroidal flow: fixed) + \mathbf{v}_1 (MHD flow)

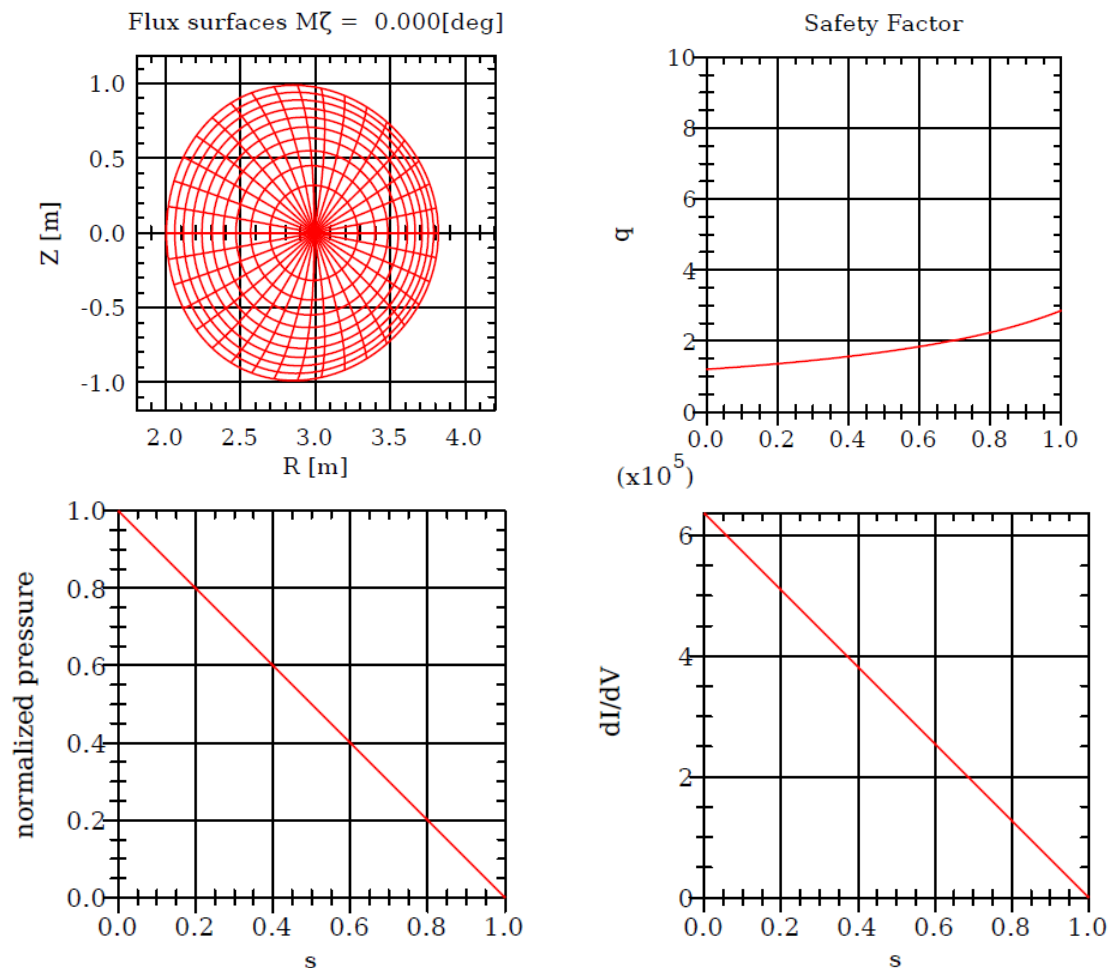
η : resistivity (constant) ν : viscosity (constant)

Collaboration with C. Hegna (U. Wisconsin)



Impact of plasma rotation on 3D MHD equilibrium

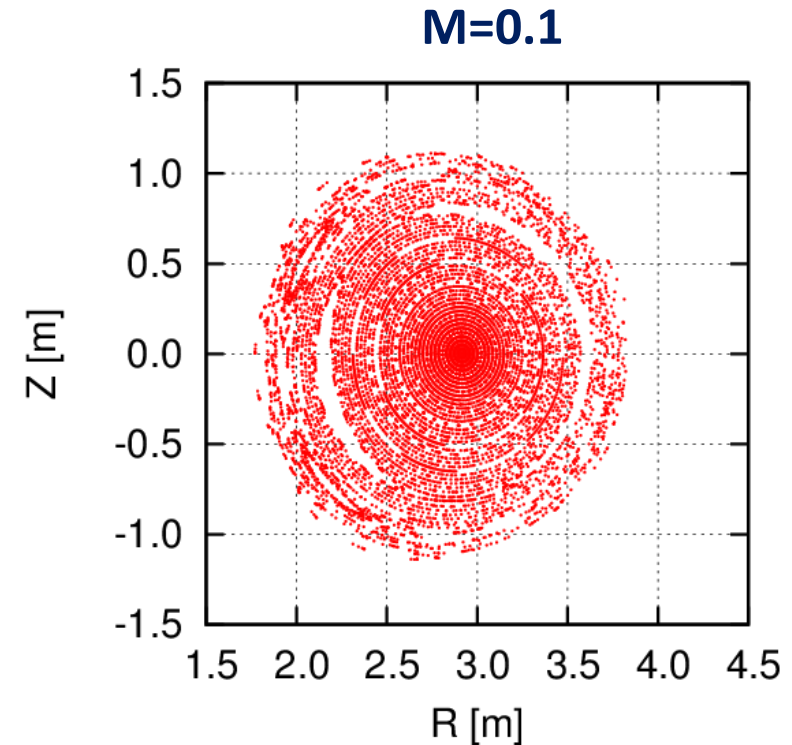
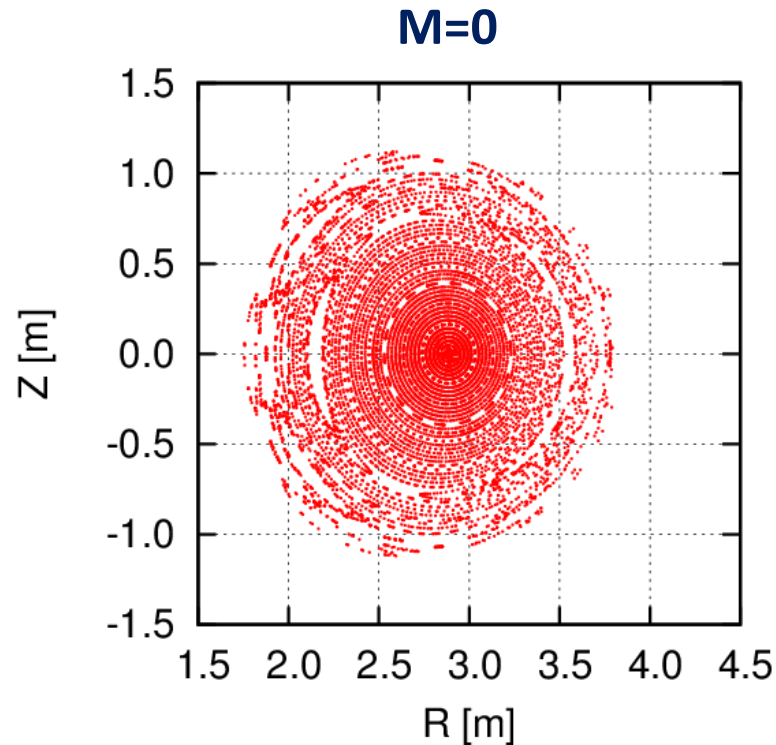
Reference equilibrium



Impact of plasma rotation on 3D MHD equilibrium

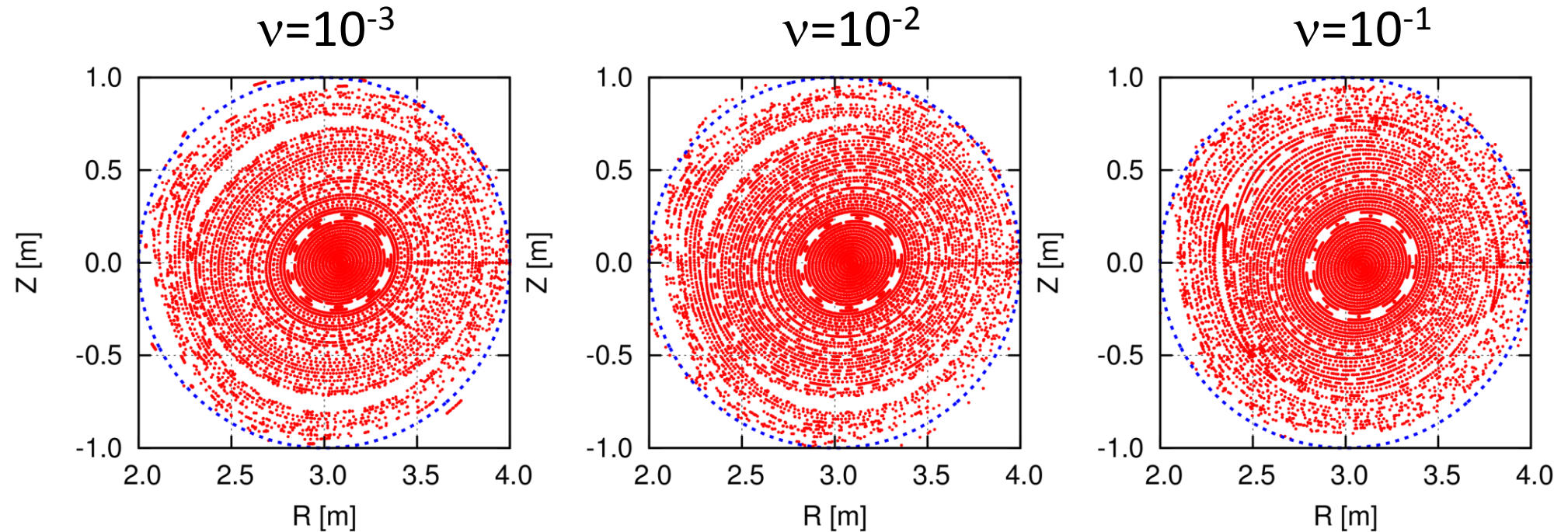
An example of improved HINT

Y.Suzuki, et al, EPS2016



- Rigid toroidal rotation is assumed.
- Phases of island slips poloidally.

Impact of perpendicular viscosity



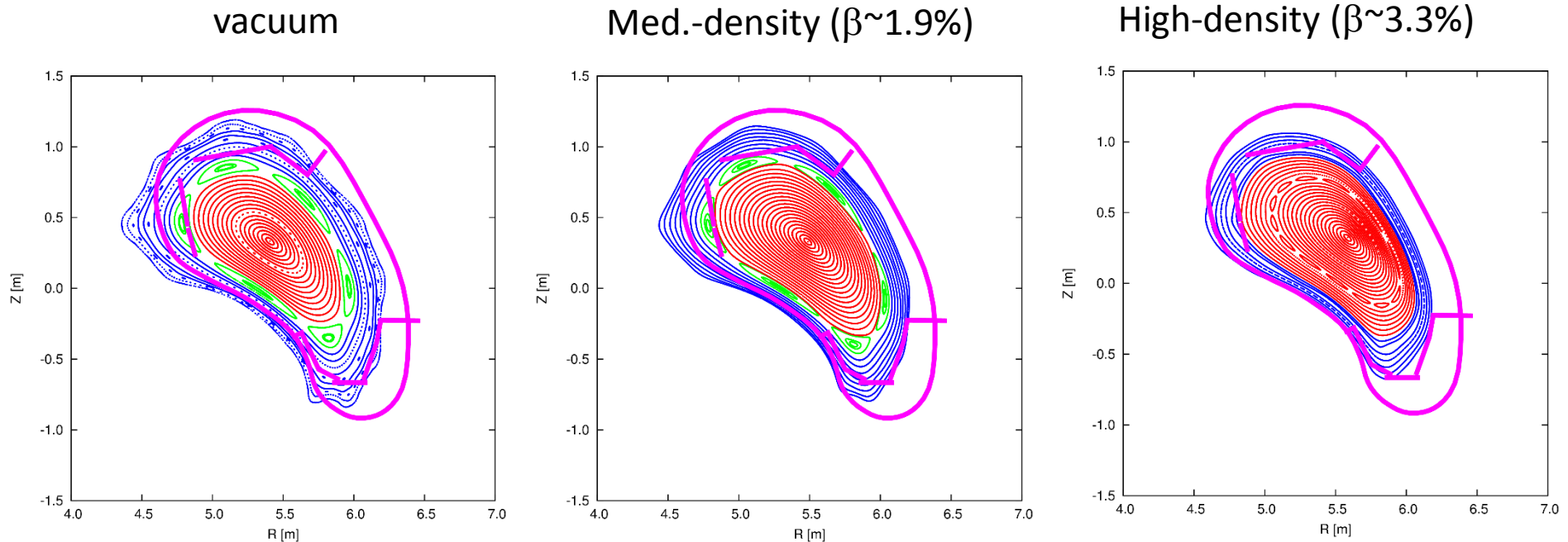
- For $\nu=10^{-3}$ and 10^{-2} , island width and phase are almost identical.
- For $\nu=10^{-1}$, island slips poloidally, phase changes π . But, **island healing is healed.**

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Summary and outlook

- **3D magnetic topology in stellarator and tokamak is discussed.**
- **To model 3D magnetic topology, 3D equilibrium is very important. Numerical codes to calculate 3D MHD equilibrium**
- **Recent progress of modeling 3D magnetic topology are discussed.**
- **For V&V, systematic comparisons of experiments and simulations are very important.**

Global effects or Local effects?



- Clear flux surfaces are kept with increasing β .
- 5/6 islands shrink due to increasing β .
- For $\beta \sim 2\%$, islands are almost healed.
- For $\beta > 3\%$, the phase of 5/6 islands changes.

Islands are healed without resonant current on 5/6 rational surface.



Contributions of **non-resonant** components.

